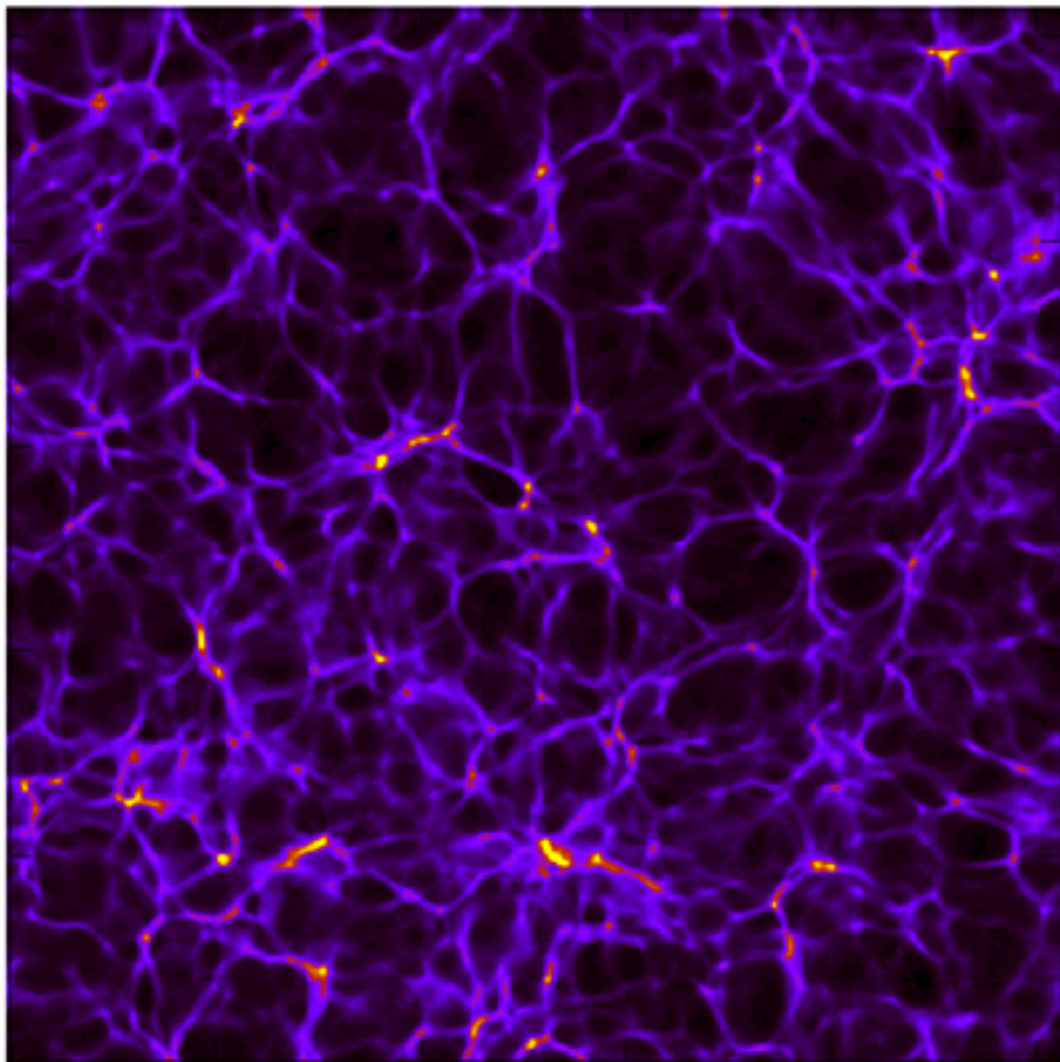


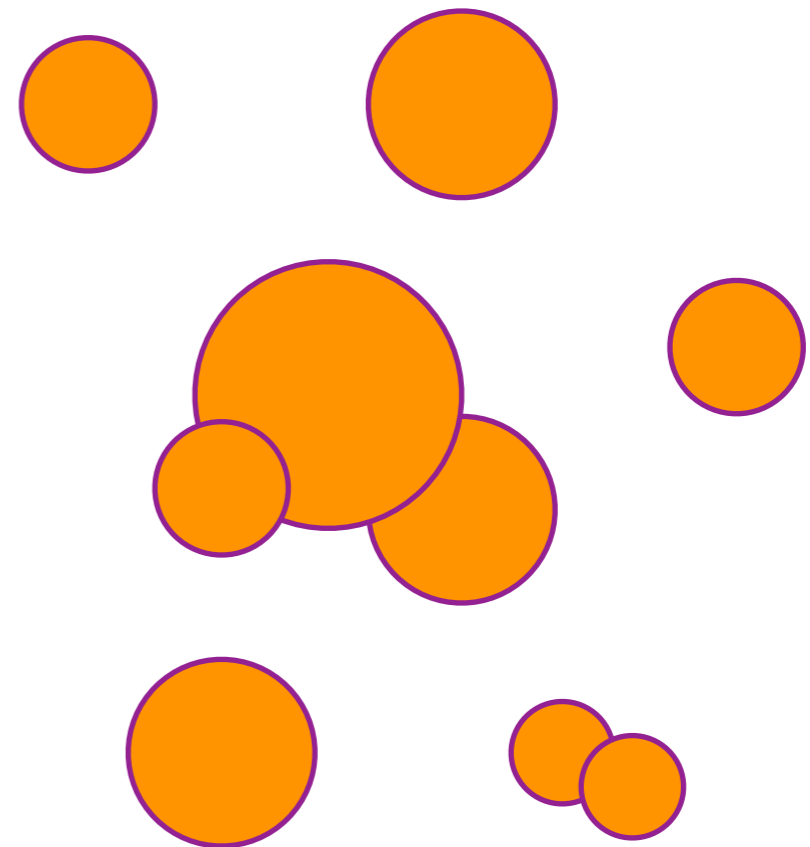
The halo model and cosmological power spectra

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Collaborators - John Peacock, Catherine Heymans, Shahab Joudaki, Alan Heavens, Lucas Lombriser



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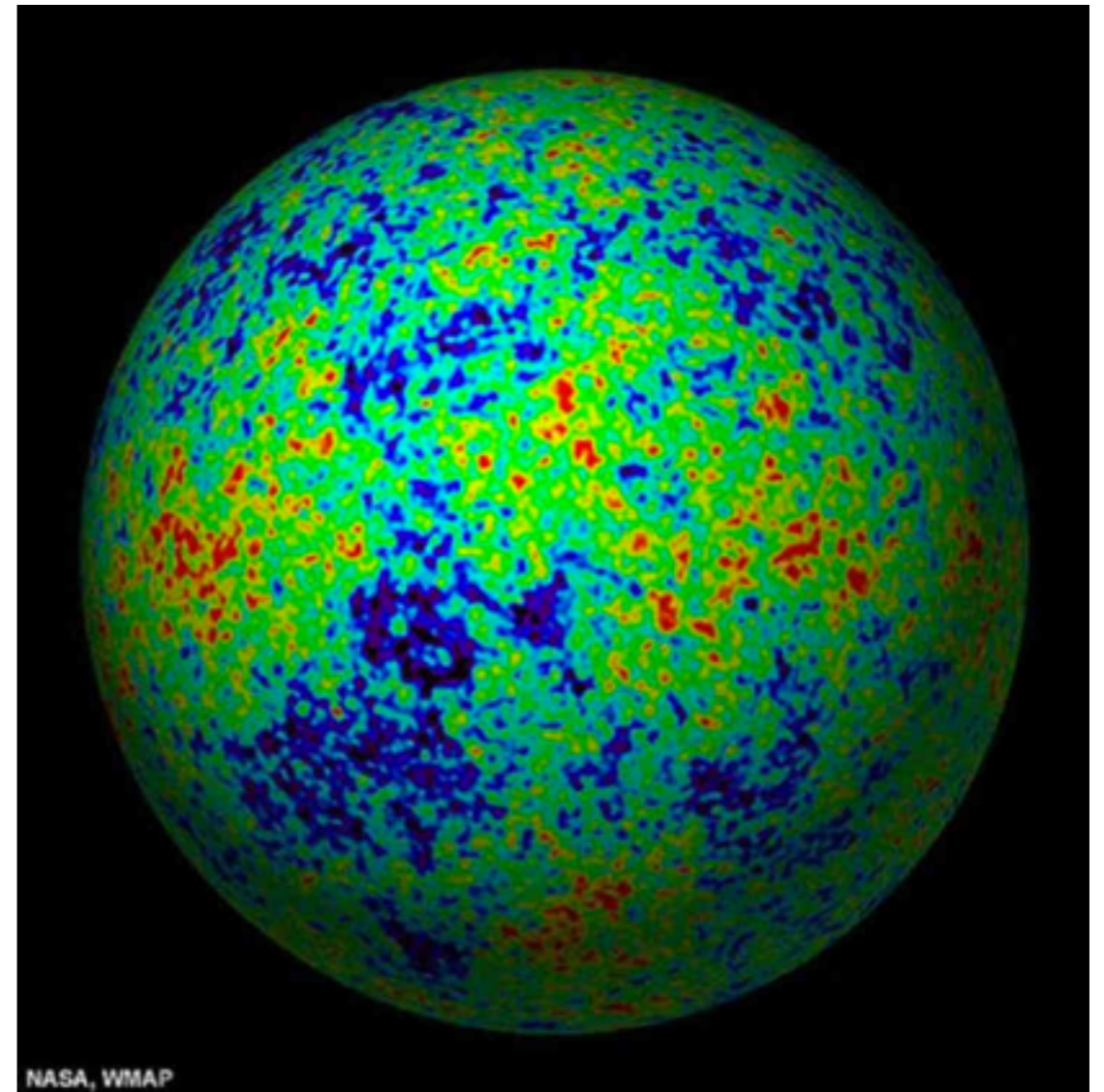


Introduction

- Aim to understand the distribution of large-scale structure (matter) in the late ($z < 2$) Universe.
- Structure is traced out by galaxies, which we can see, but consists of a dark matter exoskeleton, often called the 'cosmic web'
- This structure can be simulated easily, and measured directly via weak gravitational lensing
- Different cosmological models predict very different patterns for the cosmic web, and can thus be distinguished using weak lensing.

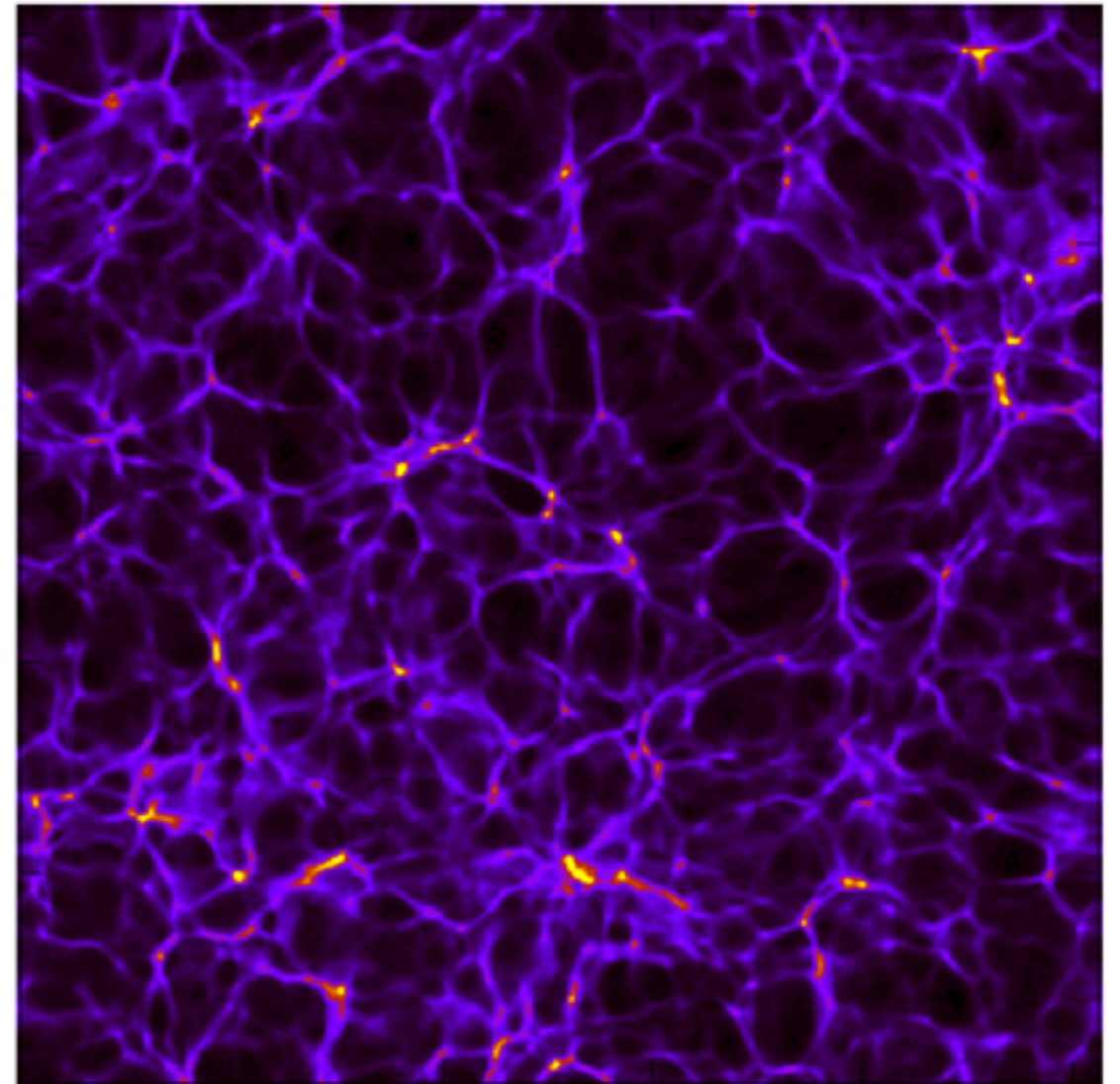
The problem

- Structure in the late Universe is non-linear and complex
- Initial perturbations are linear and Gaussian
- Gravitational collapse makes things non-linear and non-Gaussian
- Late Universe must be modelled (e.g., weak lensing)
- Necessary to do this as a function of model (DE, MG, massive neutrinos)



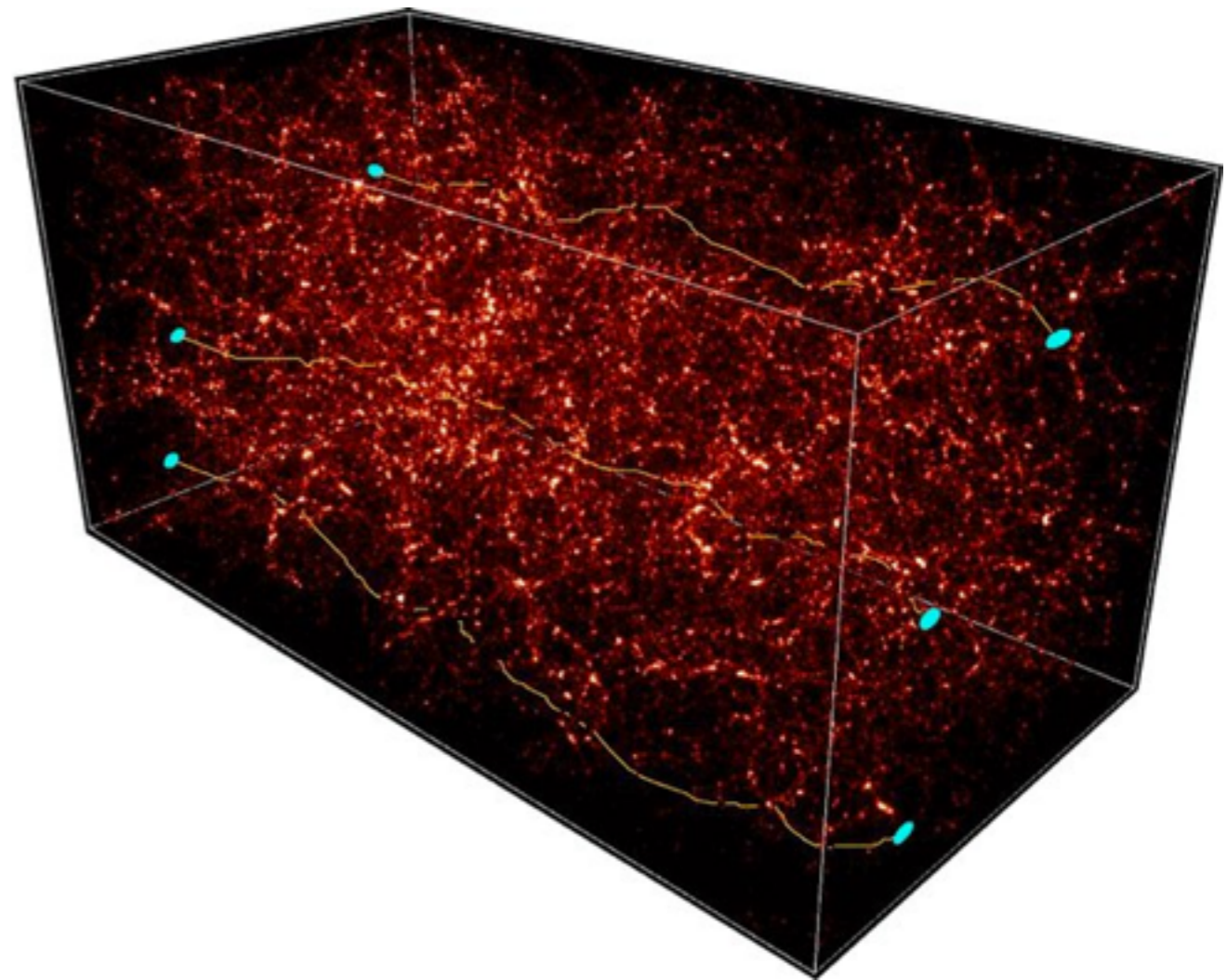
The problem

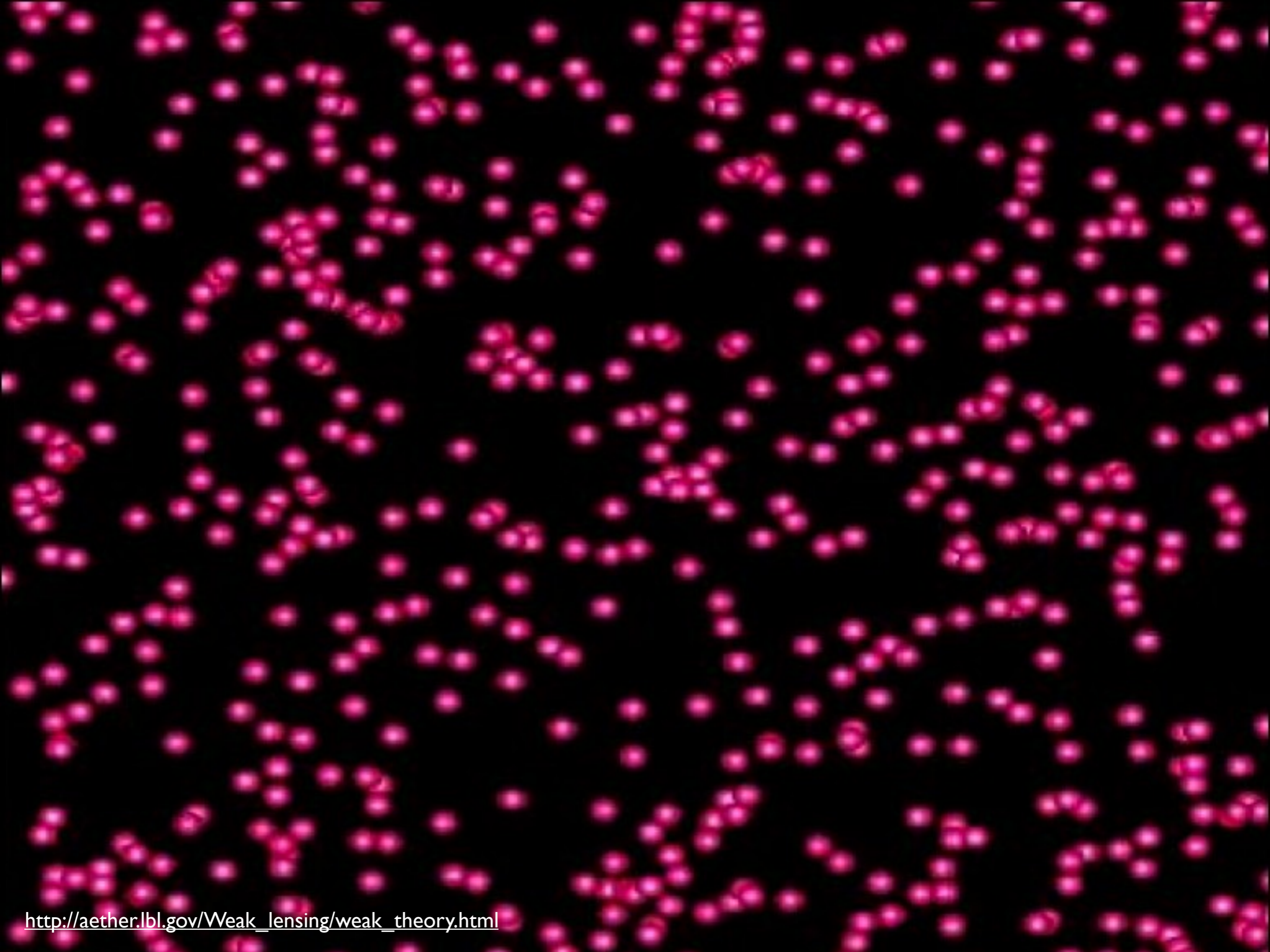
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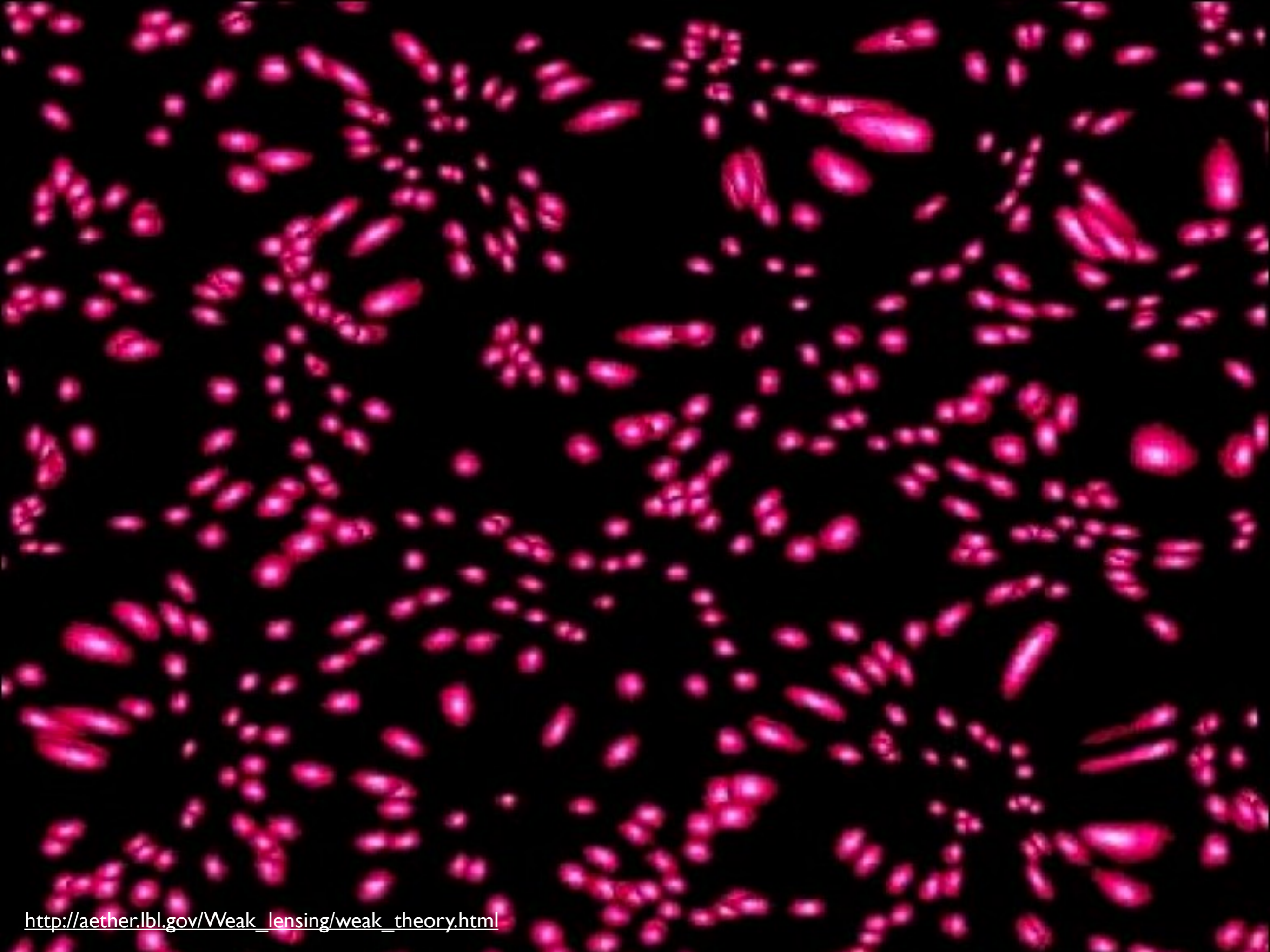


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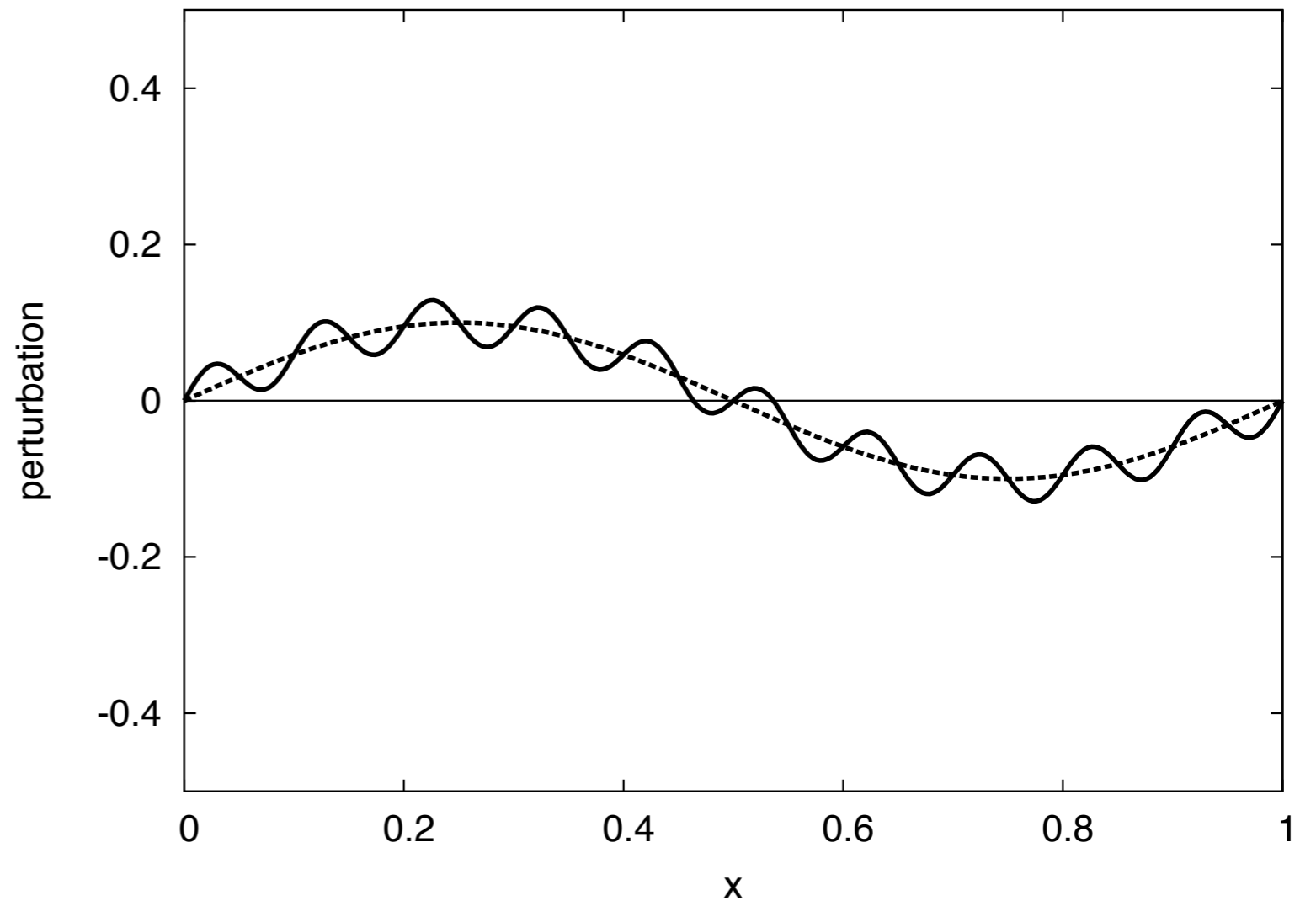




Linear perturbation theory

$$1 + \delta(\mathbf{x}) = \frac{\rho(\mathbf{x})}{\bar{\rho}}$$

$$\ddot{\delta} + 2H\dot{\delta} = 4\pi G\bar{\rho}_m\delta$$

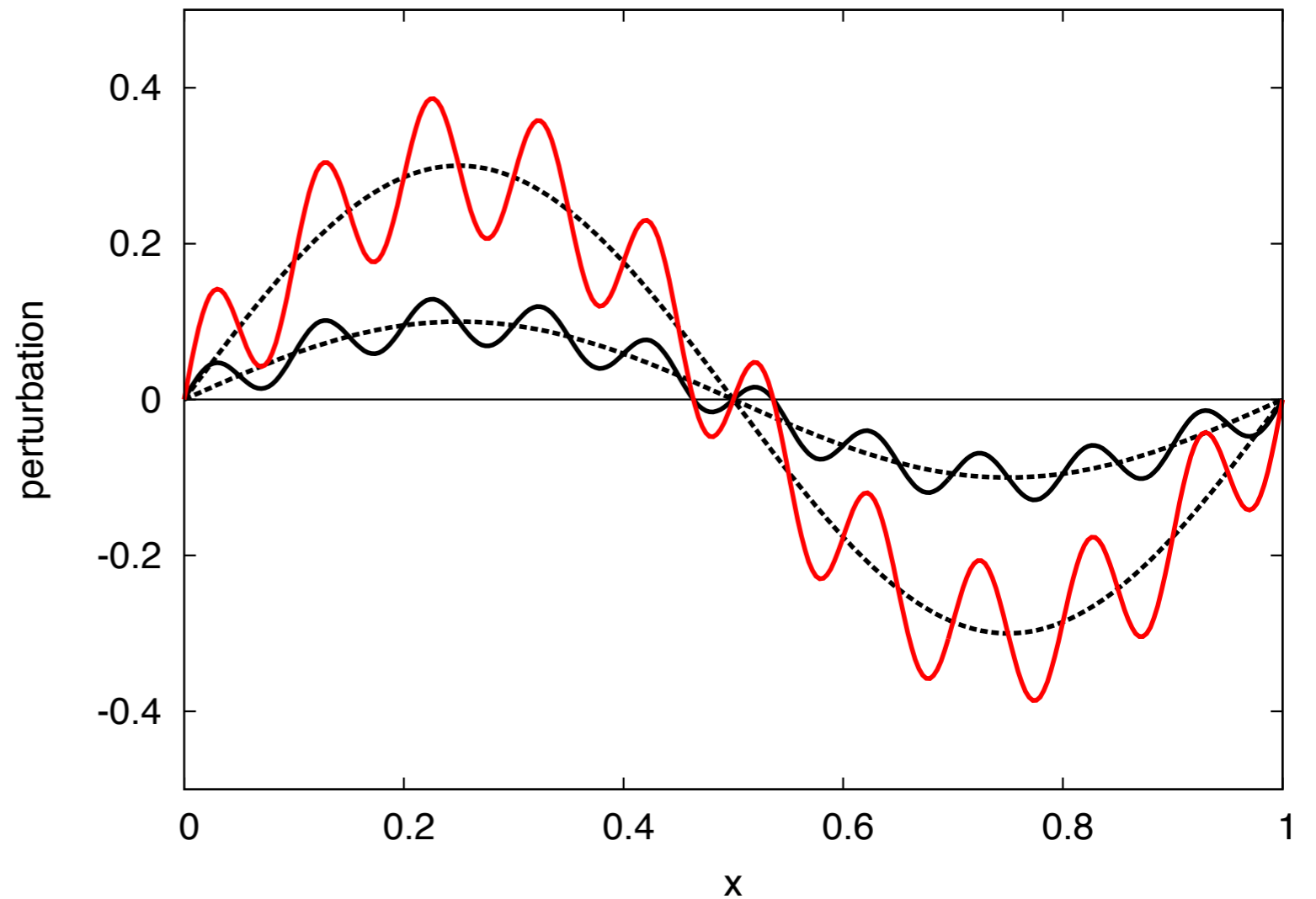


Linear perturbation theory

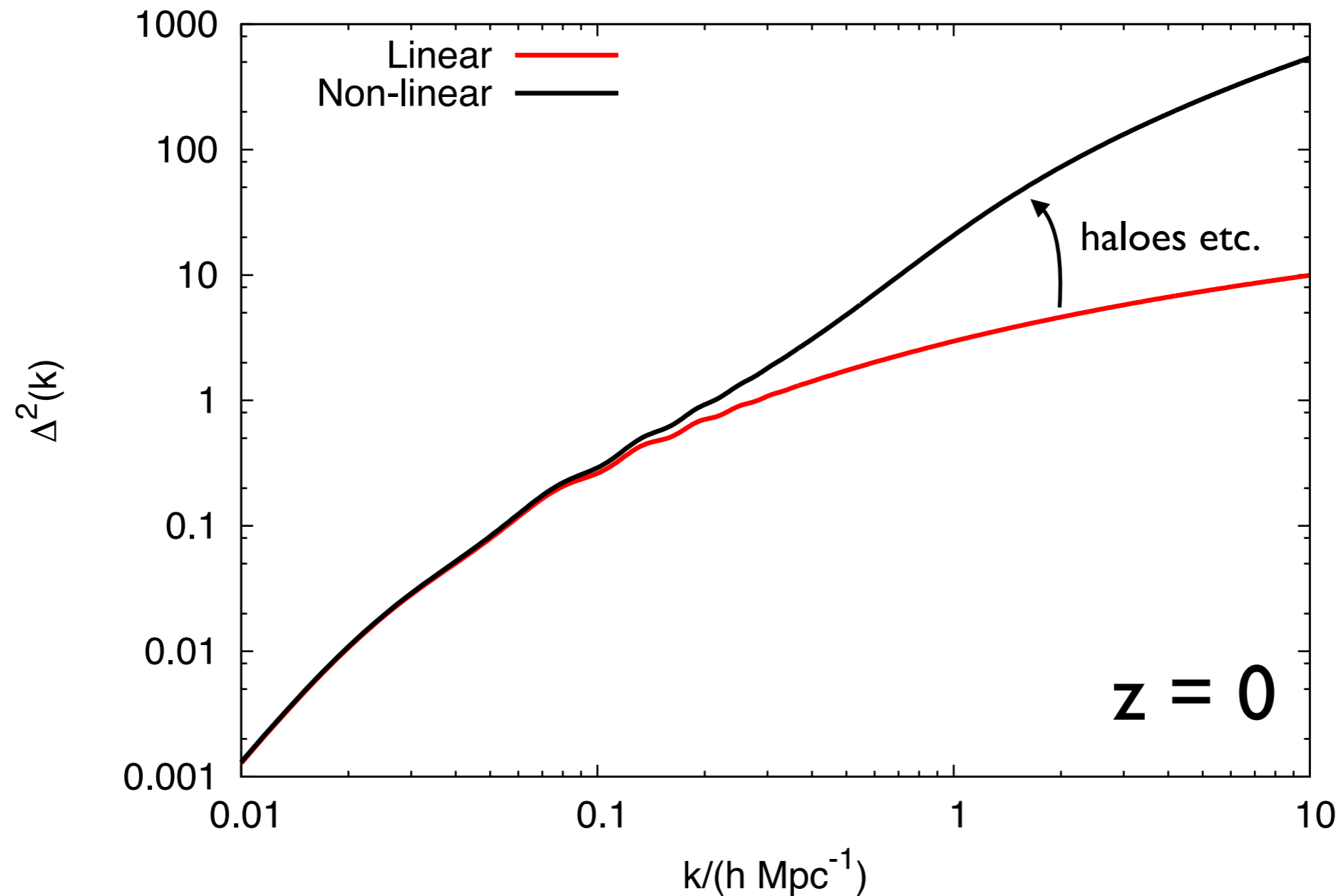
$$1 + \delta(\mathbf{x}) = \frac{\rho(\mathbf{x})}{\bar{\rho}}$$

$$\ddot{\delta} + 2H\dot{\delta} = 4\pi G\bar{\rho}_m\delta$$

$$\delta \propto a$$



The power spectrum of matter fluctuations



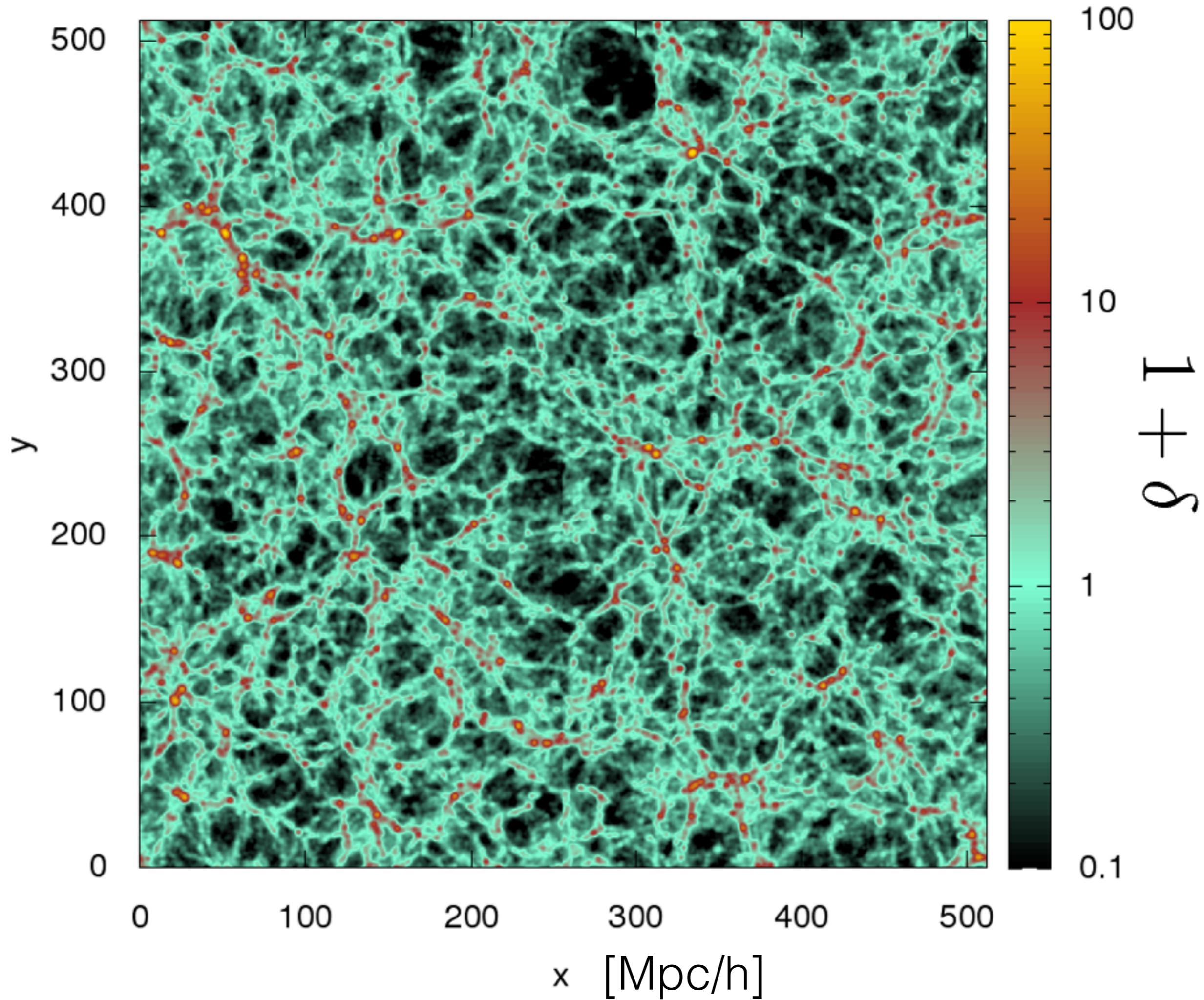
$$\Delta^2(k, z) = 4\pi V \left(\frac{k}{2\pi} \right)^3 P(k, z) \quad P(k, z) = \langle |\delta_k|^2 \rangle$$

N-body simulations



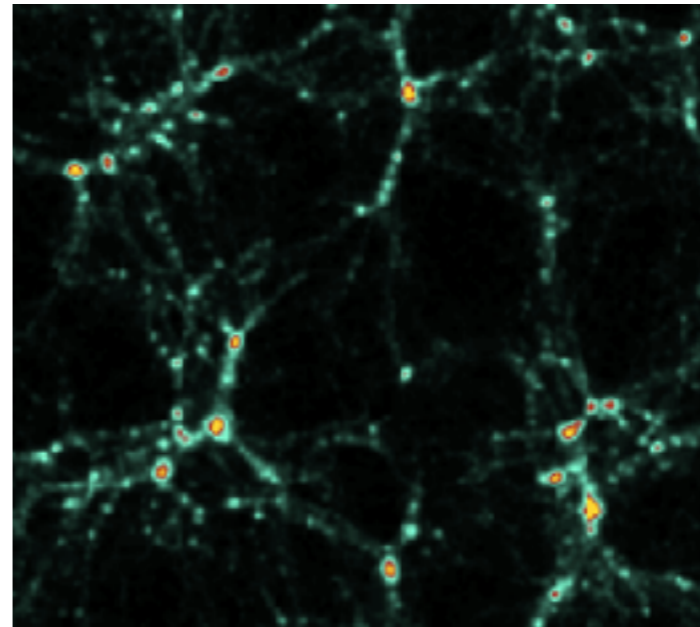
- Represent the smooth density field with 'N' particles
- 1. Calculate the gravitational forces
- 2. Move the particles
- 3. Repeat
- Periodic boundaries
- Expanding box

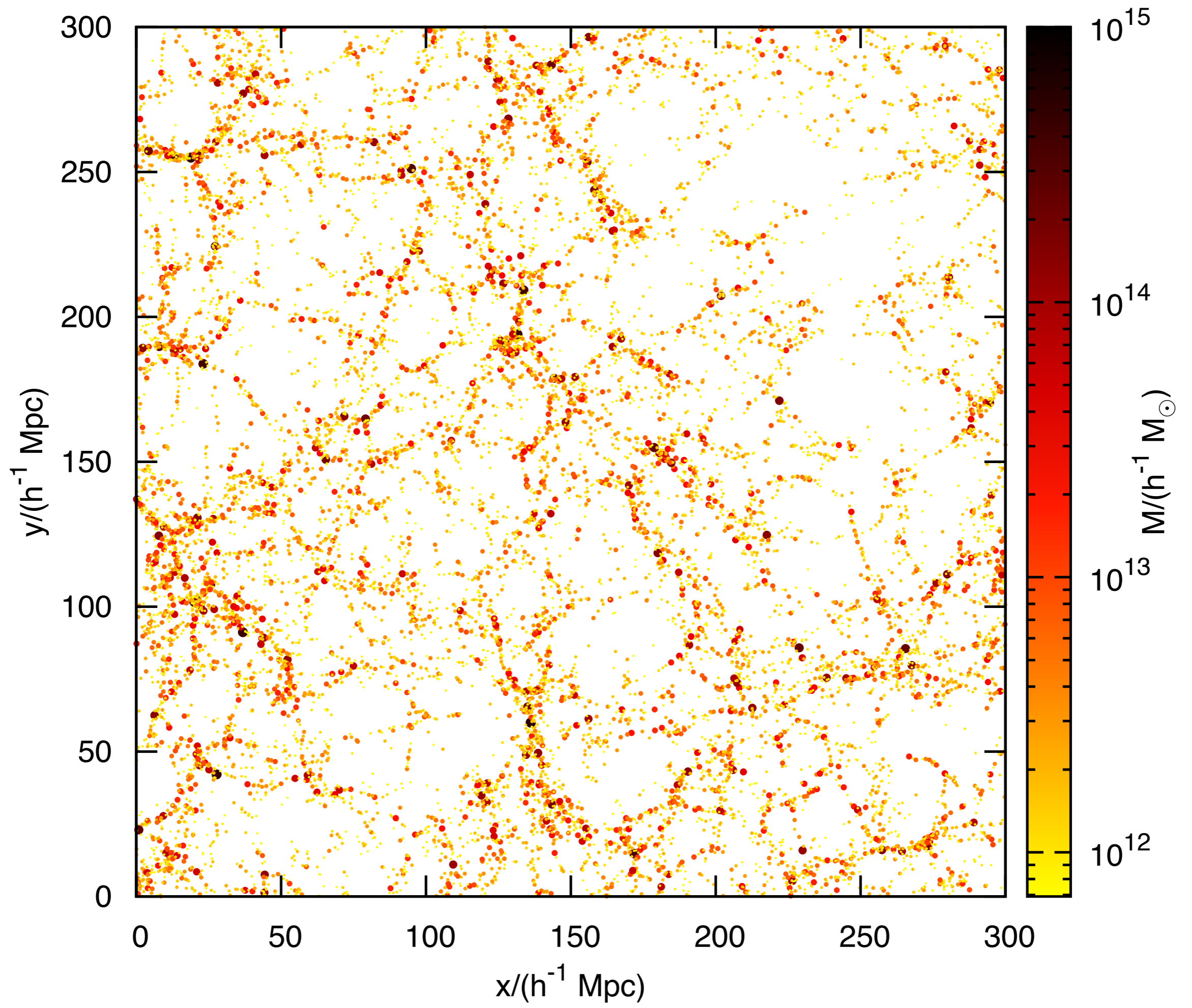
$$\ddot{\mathbf{r}} + 2H\dot{\mathbf{r}} = -\frac{1}{a^2}\nabla\Phi$$



Simulation limitations

- Solve gravity only problem to high accuracy
- Expensive in terms of computing power
- It will **never** be possible to run accurate simulations for every model under consideration
- Investigating cosmological parameter space:
 - Λ CDM
 - Dark Energy (interacting?)
 - Massive neutrinos
 - Dark matter types
 - Modified gravity
 - Baryons?





The halo model

- Can be used to predict the clustering of matter (as well as of haloes or of galaxies) in the Universe
- Distinct from HALOFIT (Smith et al. 2003)

large-scales \sim linear

$$\Delta^2(k) = \Delta_{2H}^2 + \Delta_{1H}^2$$

small scales -
contribution
from haloes

$$\Delta_{1H}^2(k) = 4\pi \left(\frac{k}{2\pi}\right)^3 \frac{1}{\bar{\rho}^2} \int_0^\infty M^2 W^2(k, M) f(M) dM$$

$$\Delta_{1H}^2(k) = 4\pi \left(\frac{k}{2\pi}\right)^3 \frac{1}{\bar{\rho}^2} \int_0^\infty M^2 W^2(k, M) f(M) dM$$

Ingredients

$$f(v) = A \left[1 + \frac{1}{(av^2)^p} \right] e^{-av^2/2}$$

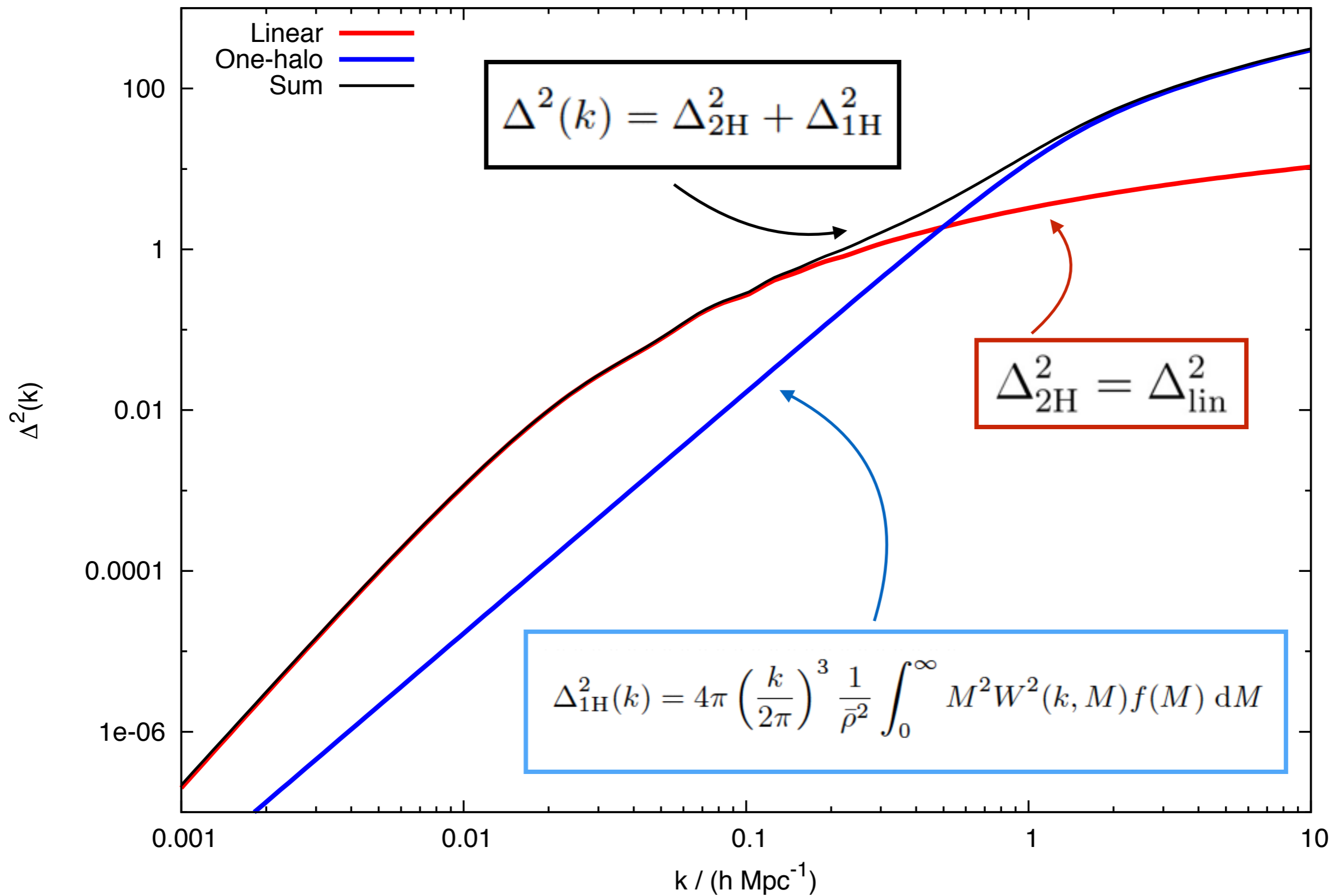
$$v = \frac{\delta_c}{\sigma(M)}$$

$$\rho(r) = \frac{\rho_N}{(r/r_s)(1+r/r_s)^2}$$

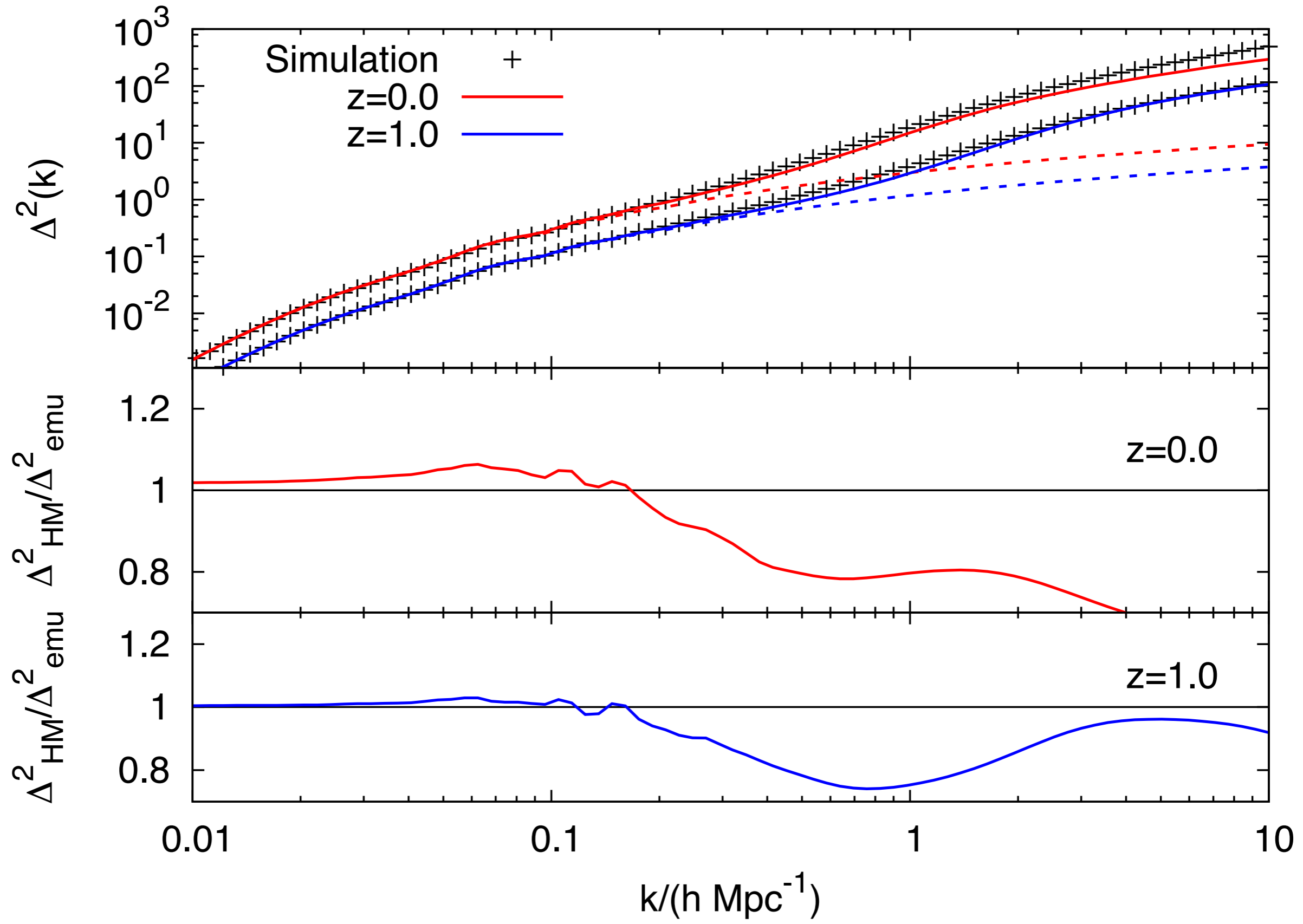
$$r_v = \left(\frac{3M}{4\pi\Delta_v\bar{\rho}} \right)^{1/3}$$

$$c(M, z) = A \frac{1+z_f}{1+z}$$

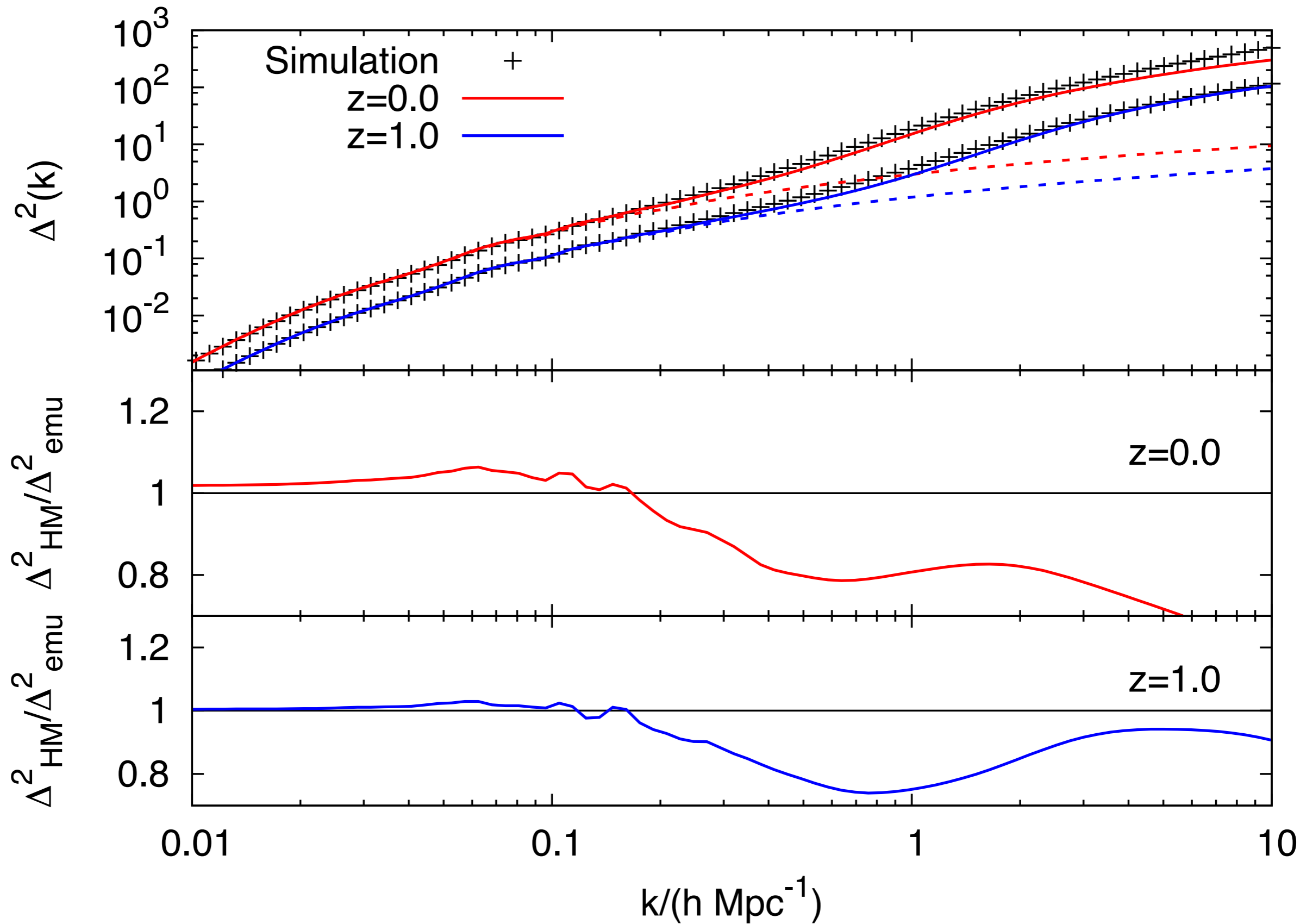
- Halo mass function
- Peak to mass relation
- Halo density profiles
- Halo radius
- Halo concentration



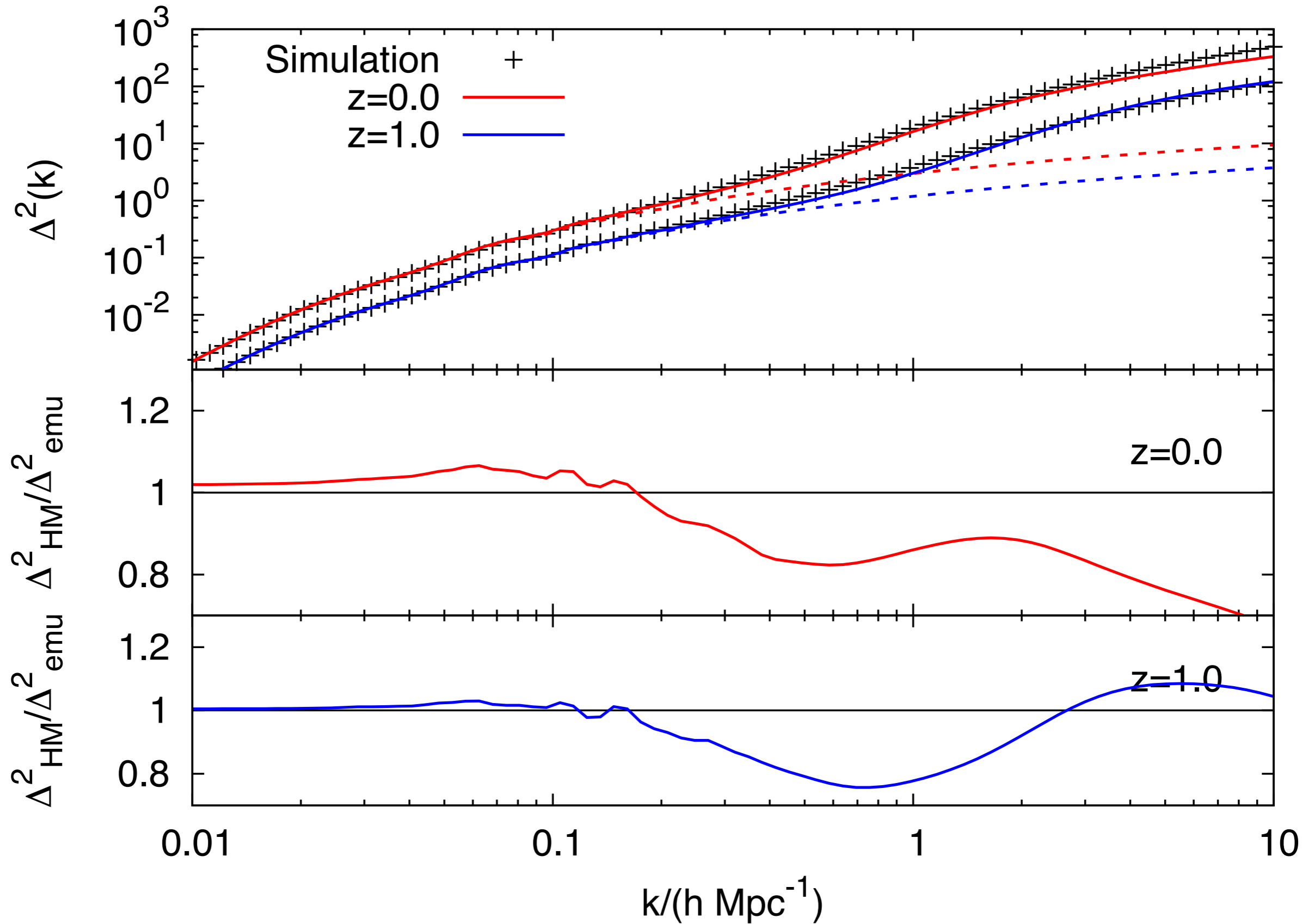
Original



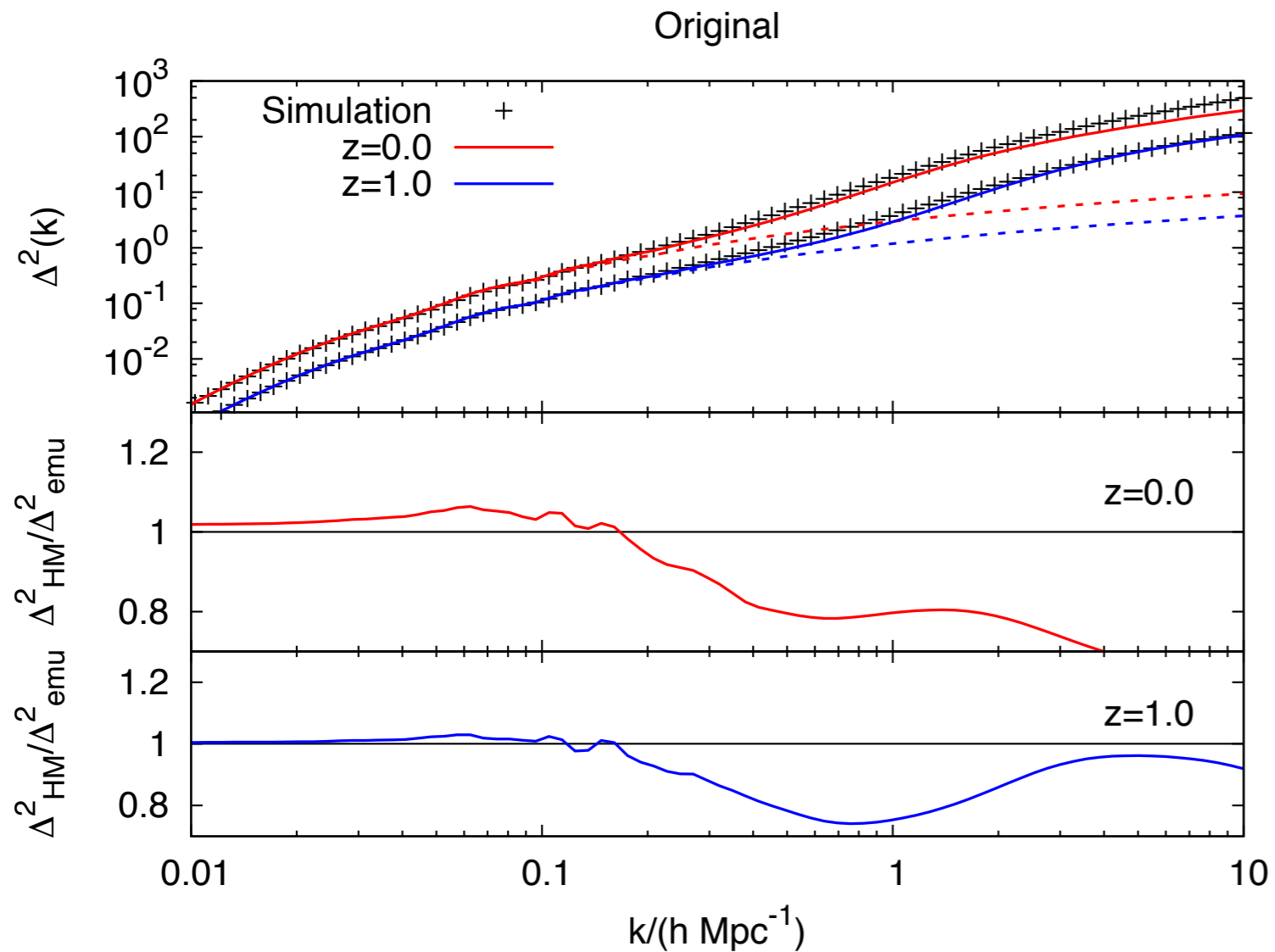
Duffy concentration-mass relation



Tinker mass function



Standard halo model

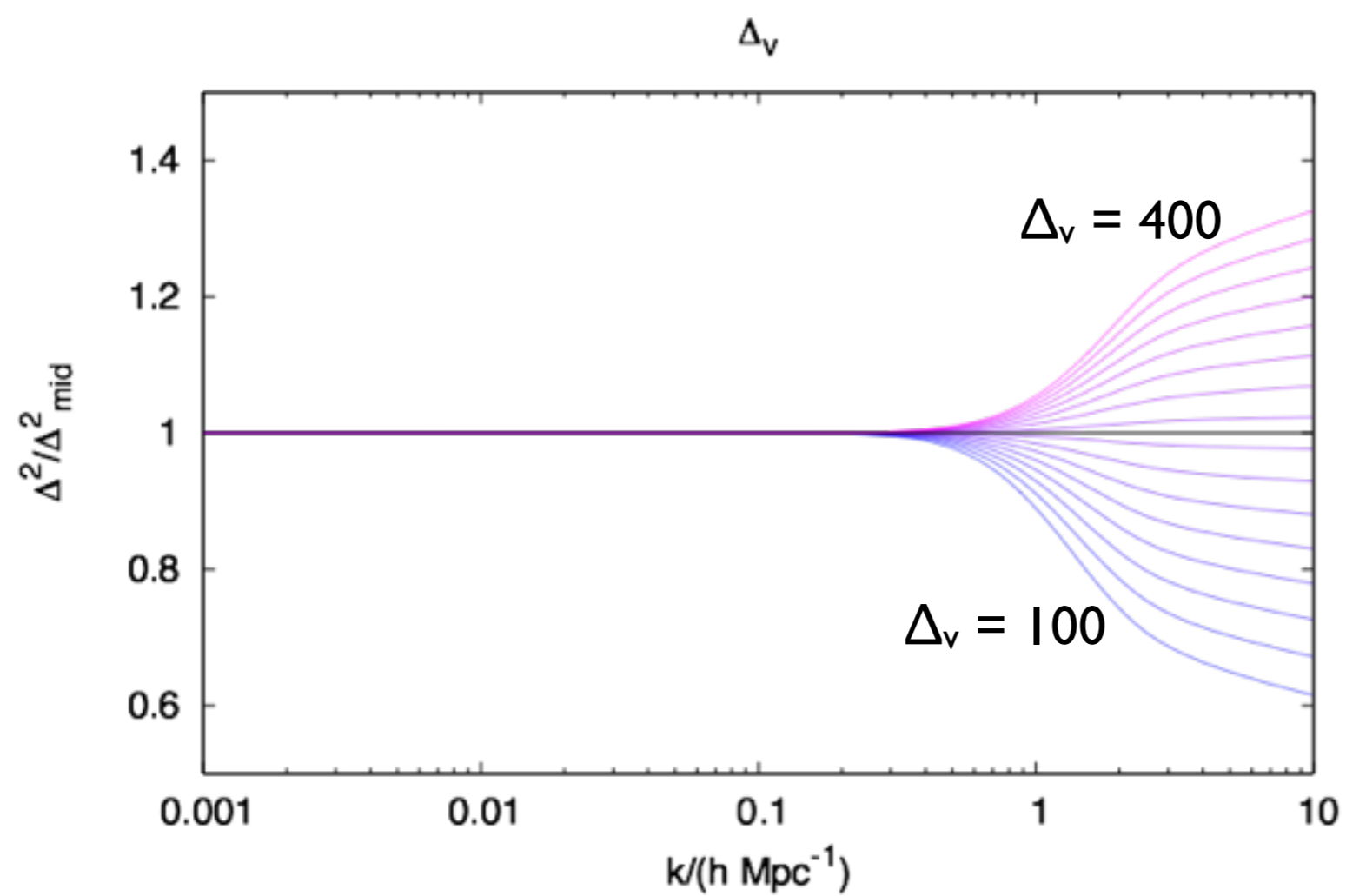


Problems

- Perturbation theory at large scales
- Transition region is problematic (voids, under-densities, NL bias)
- Filamentary structure missing
- Halo asphericity ignored
- Tidal alignment of haloes
- Assumes all objects virialized
- Halo substructure ignored
- Scatter in halo properties at fixed mass ignored

Method

- We want to remedy the inaccuracy of the halo model without breaking it
- Fixing all of the problems whilst being fully consistent would be hard
- Instead opt for the simpler goal of generating 'effective haloes', whose simple halo power spectrum accurately matches that of simulations

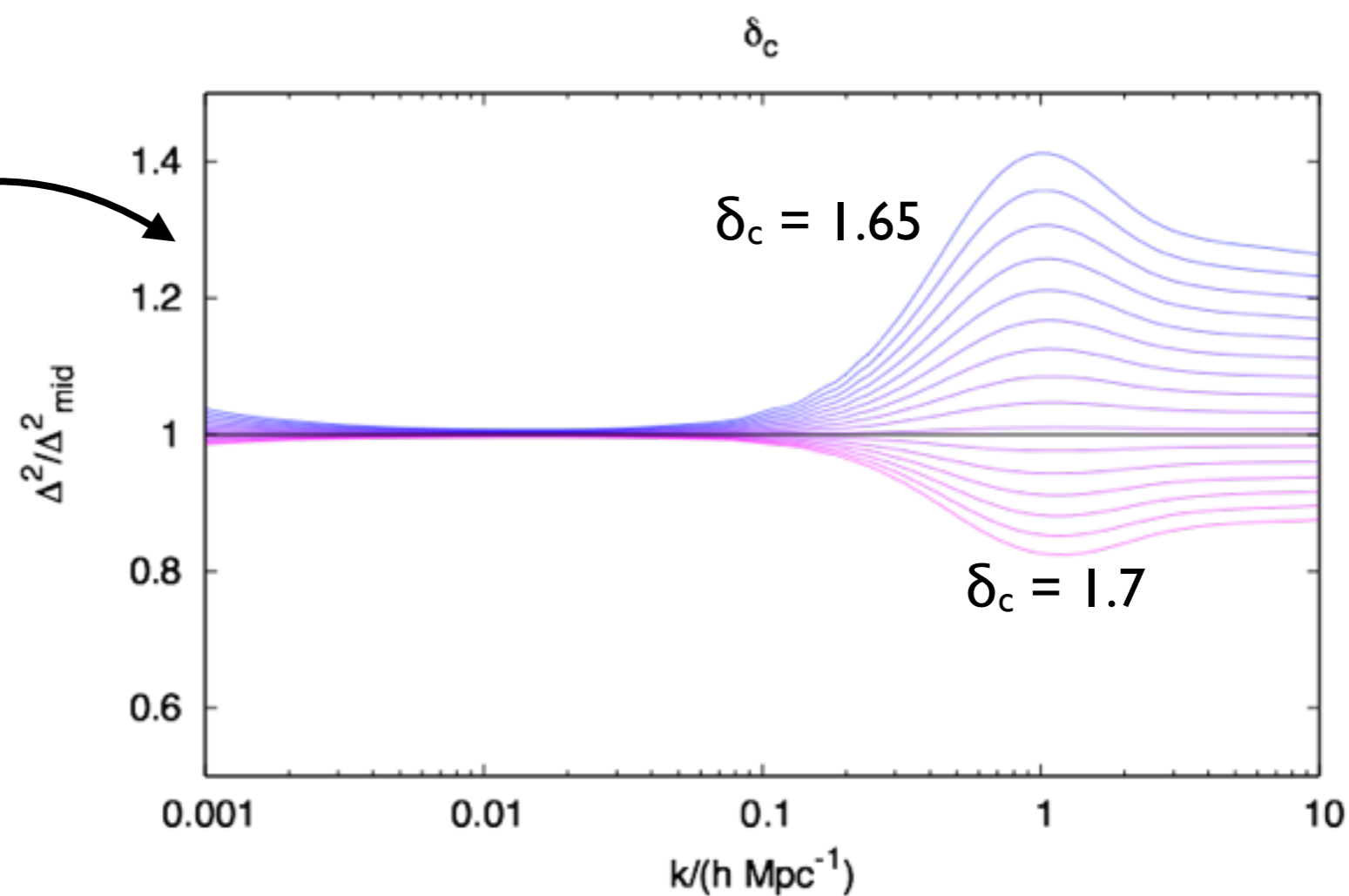


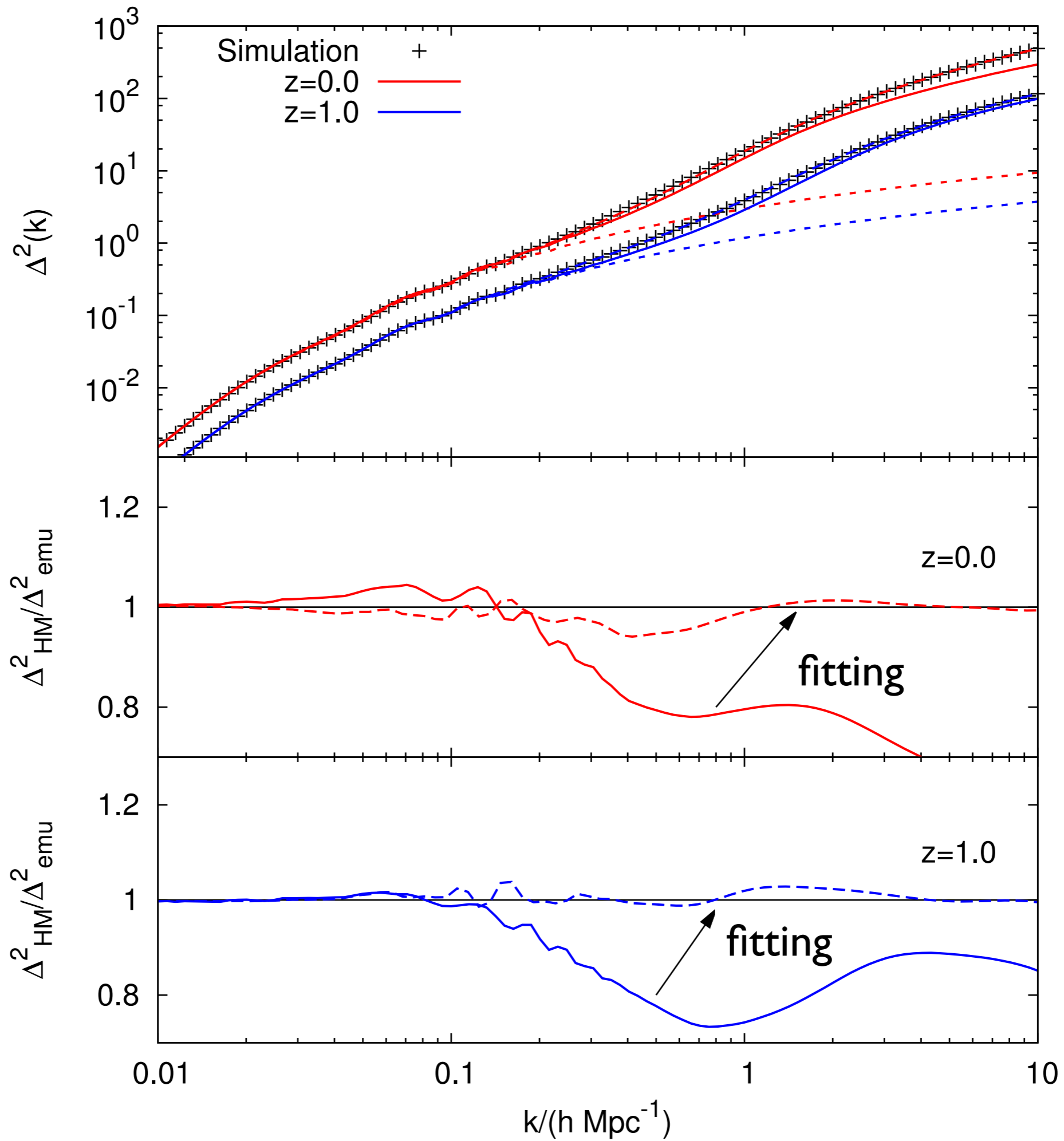
Variations in
virial density

$\Delta_v = 200$

Variations in
linear collapse
density

$\delta_c = 1.686$

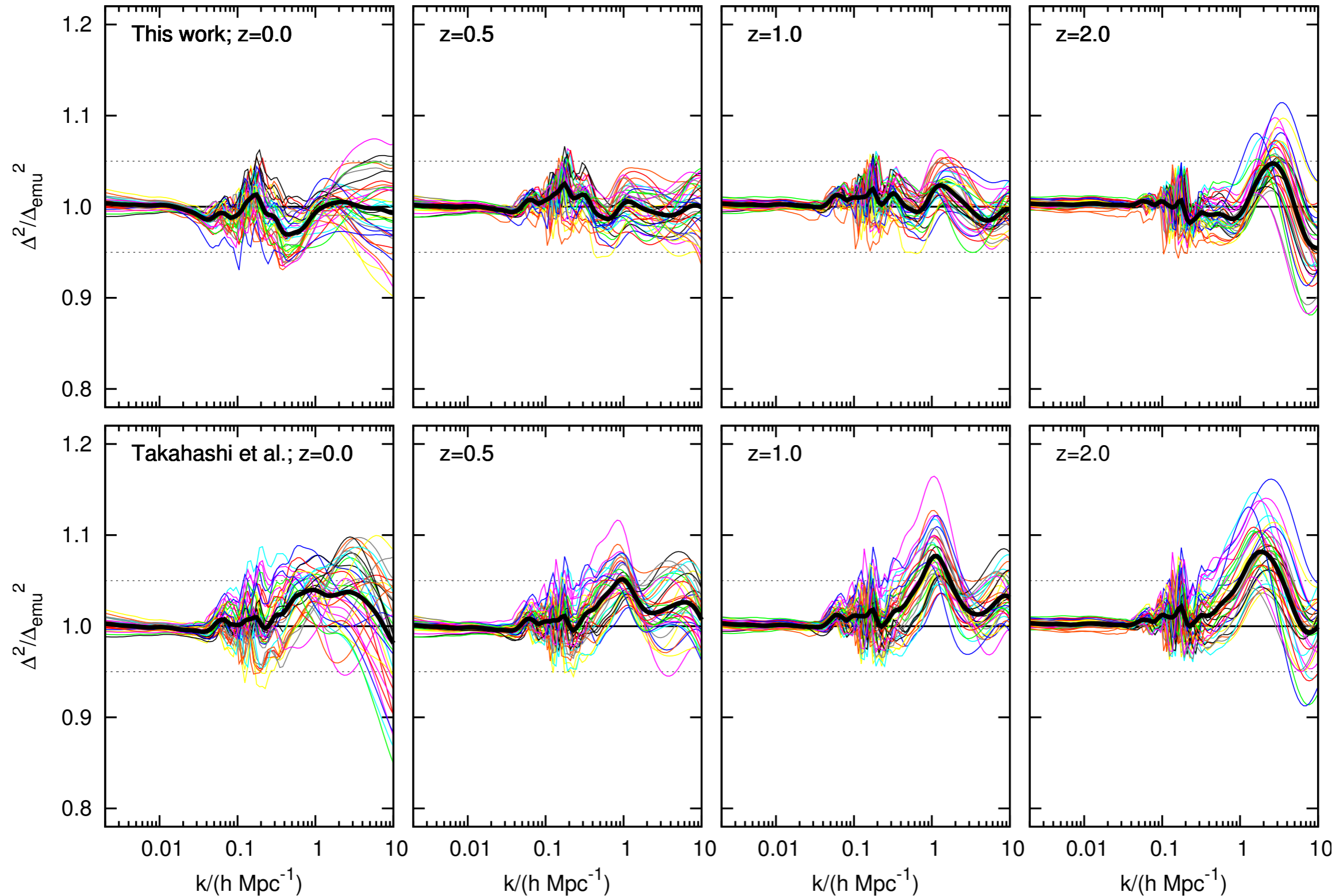




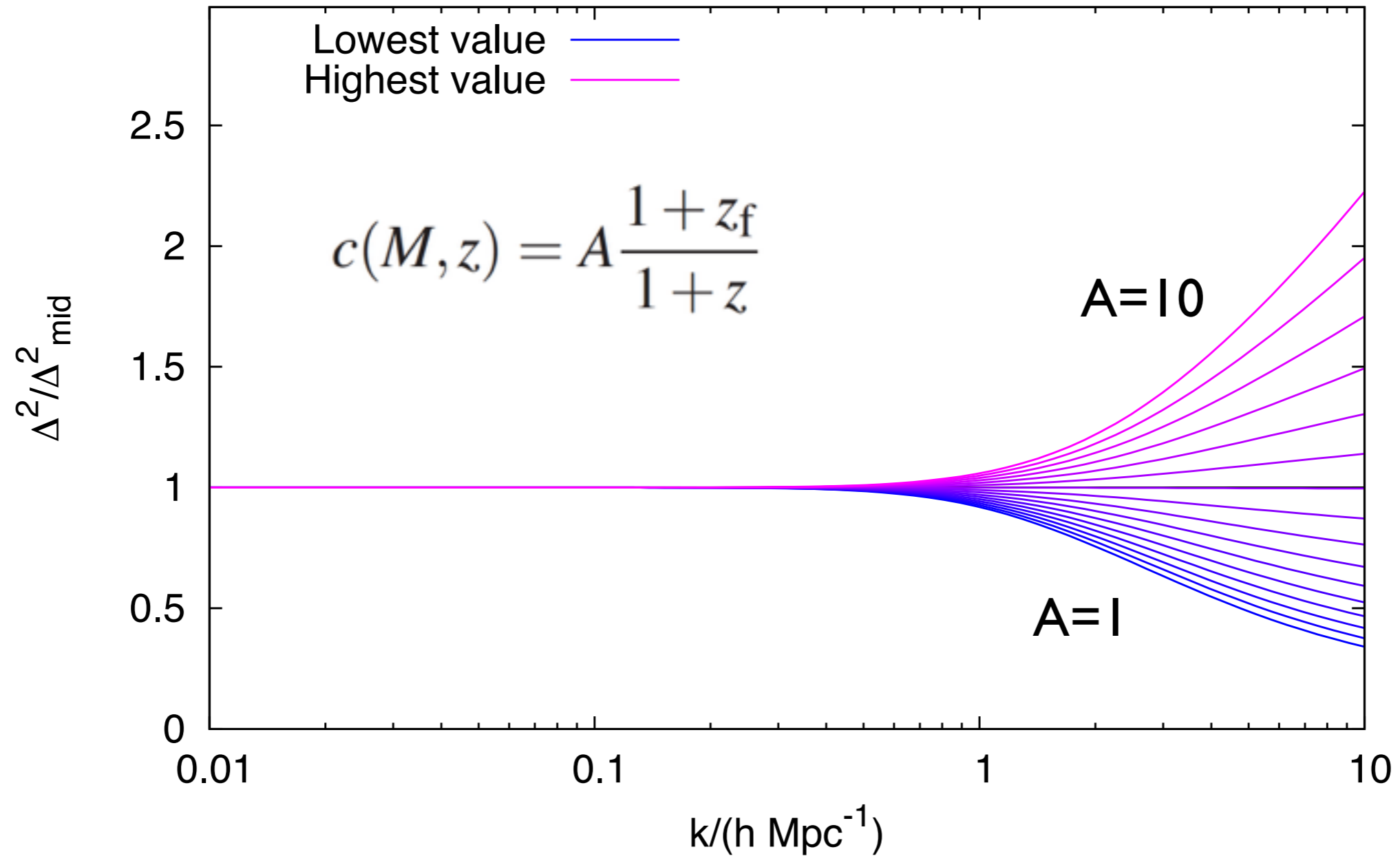
Cosmic Emu (simulations)

fit and comparison

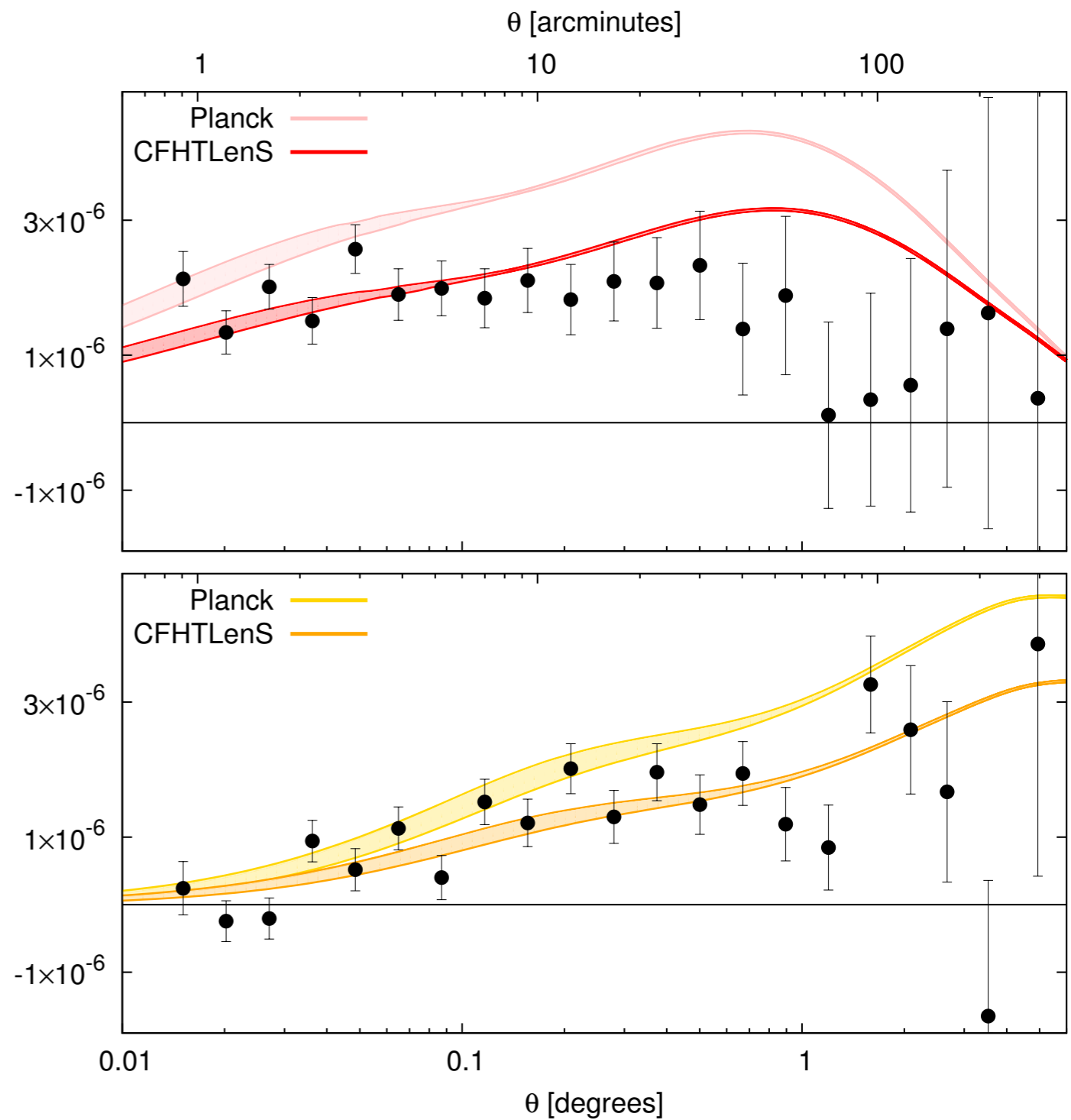
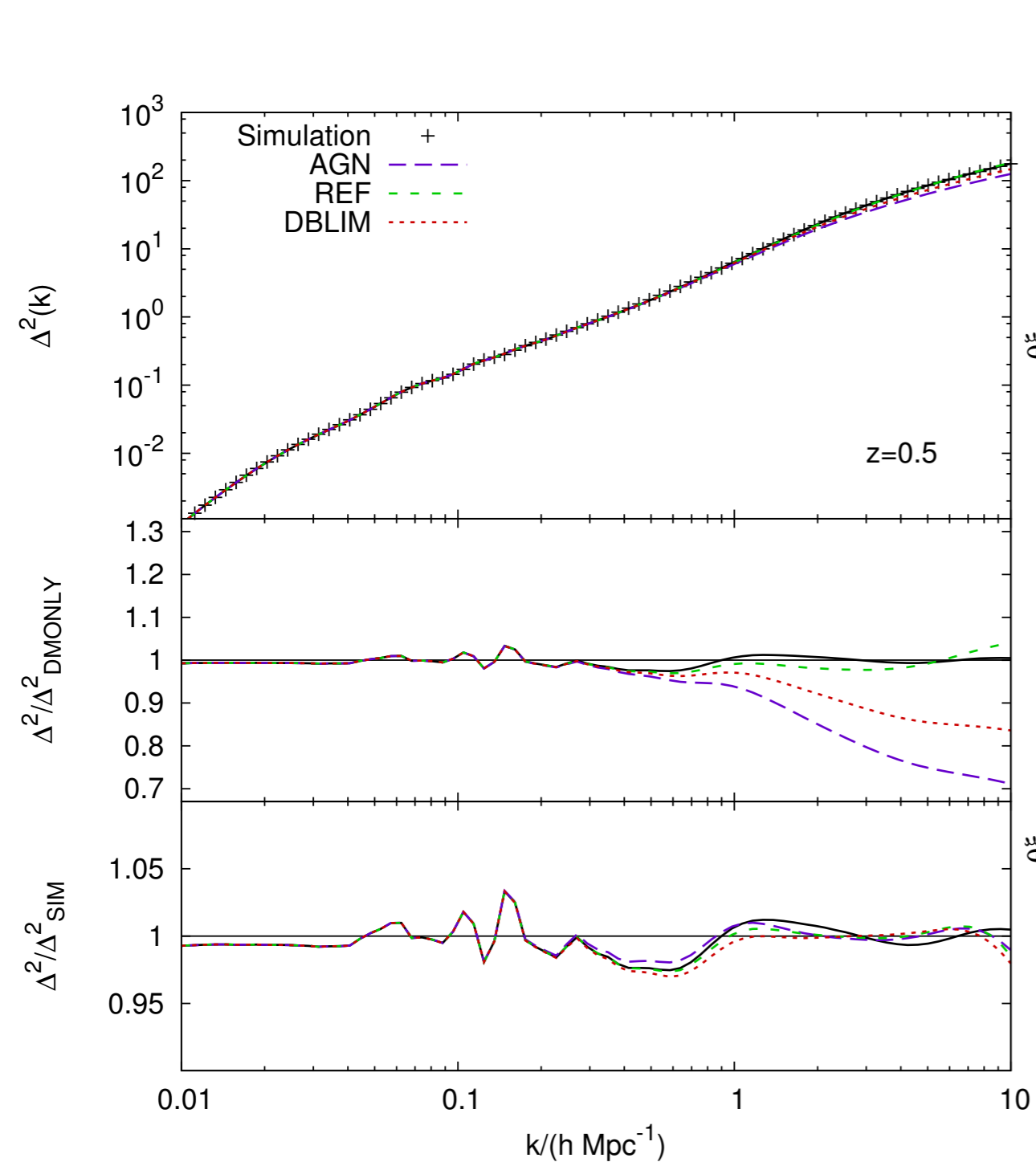
(Heitmann 2014)



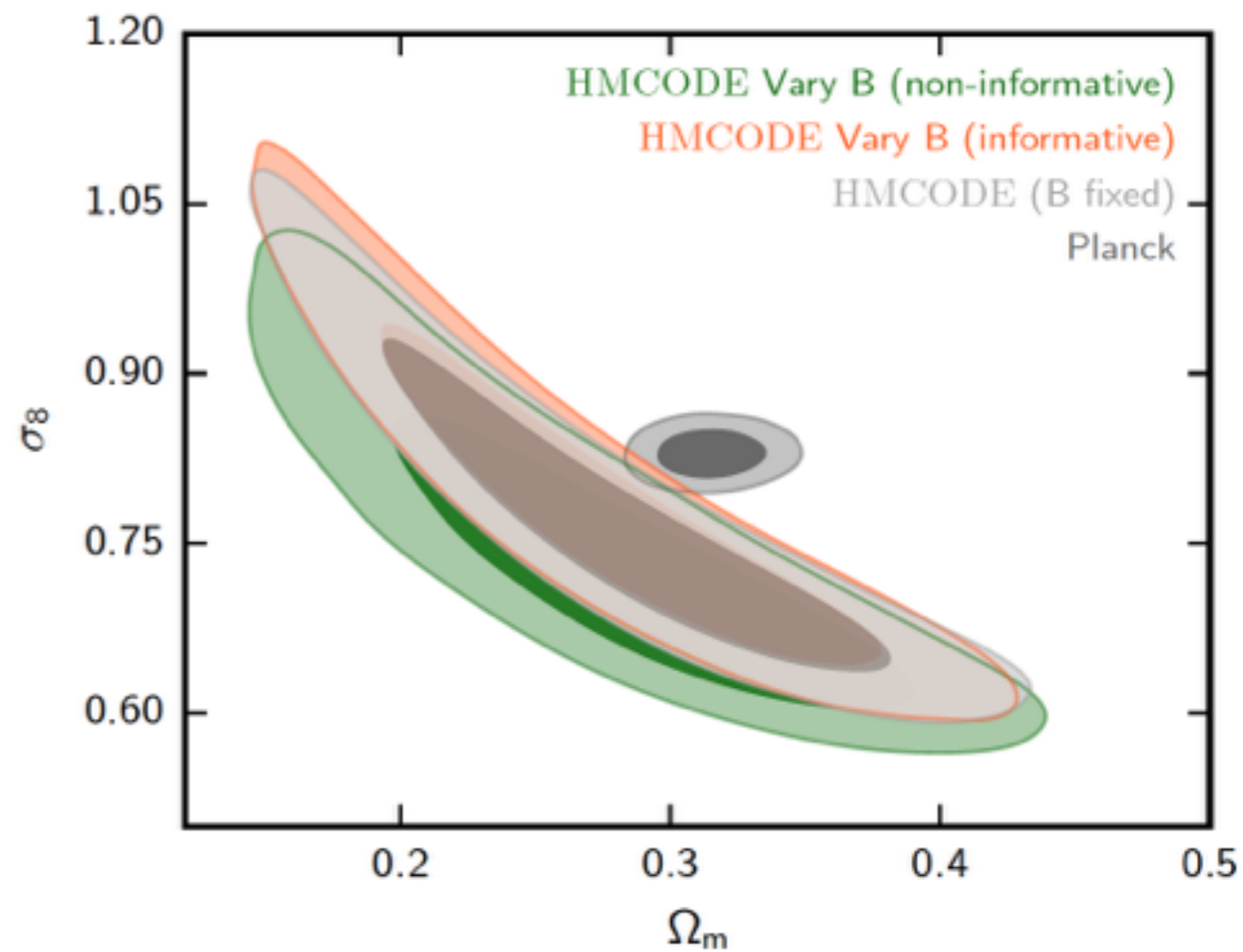
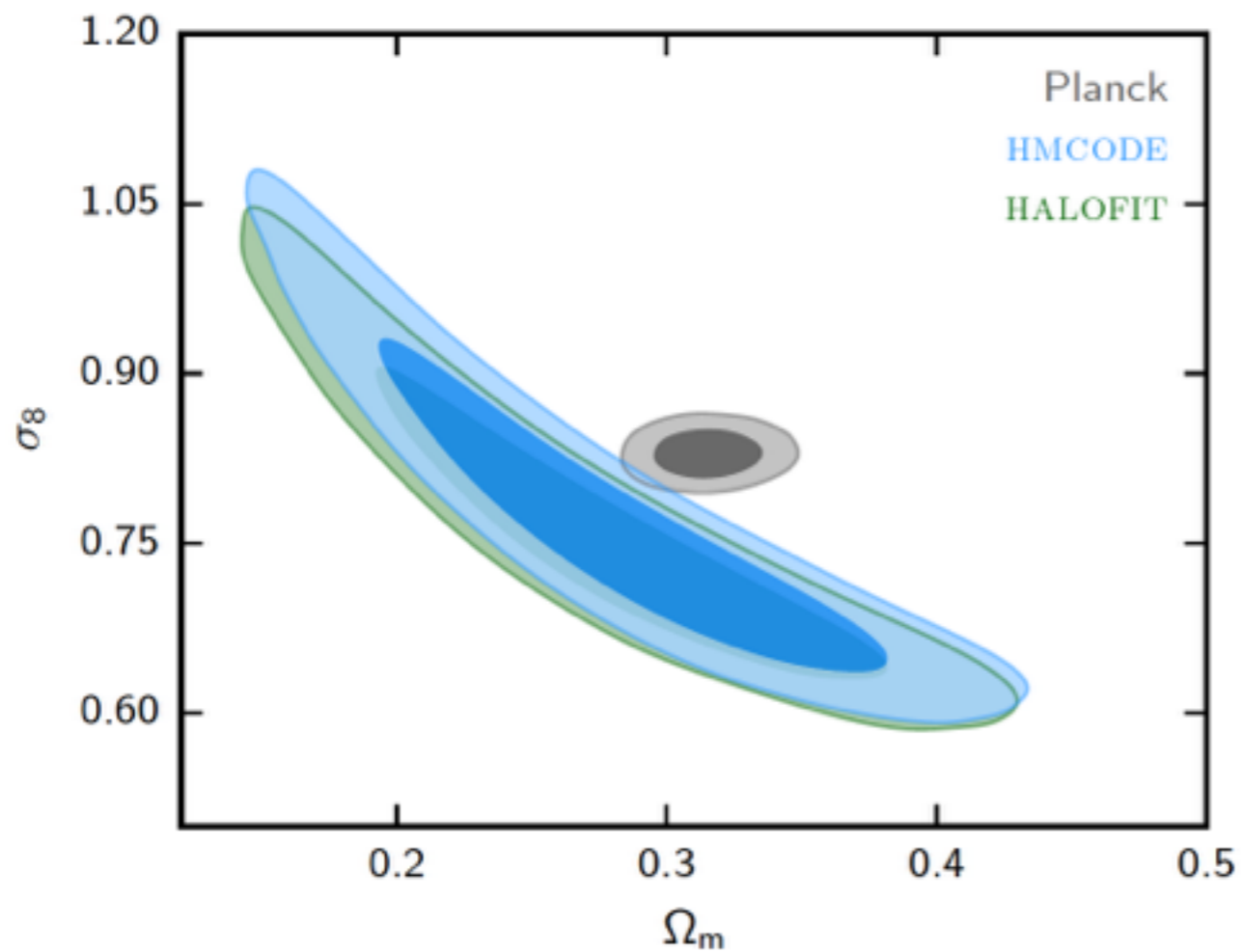
Baryonic feedback



Baryonic feedback

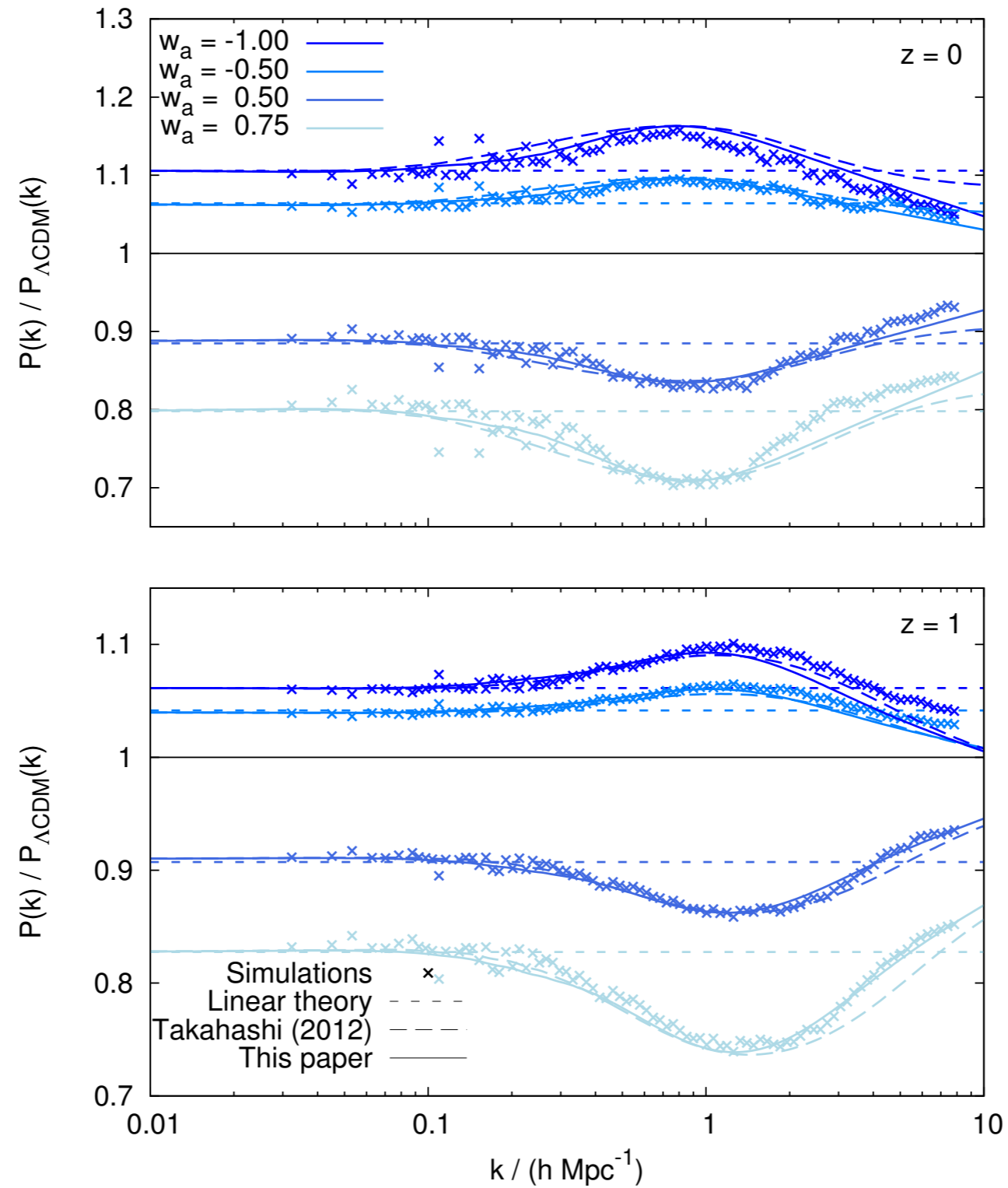


From Joudaki et al. (CFHTLenS revisited; 1601.05786)

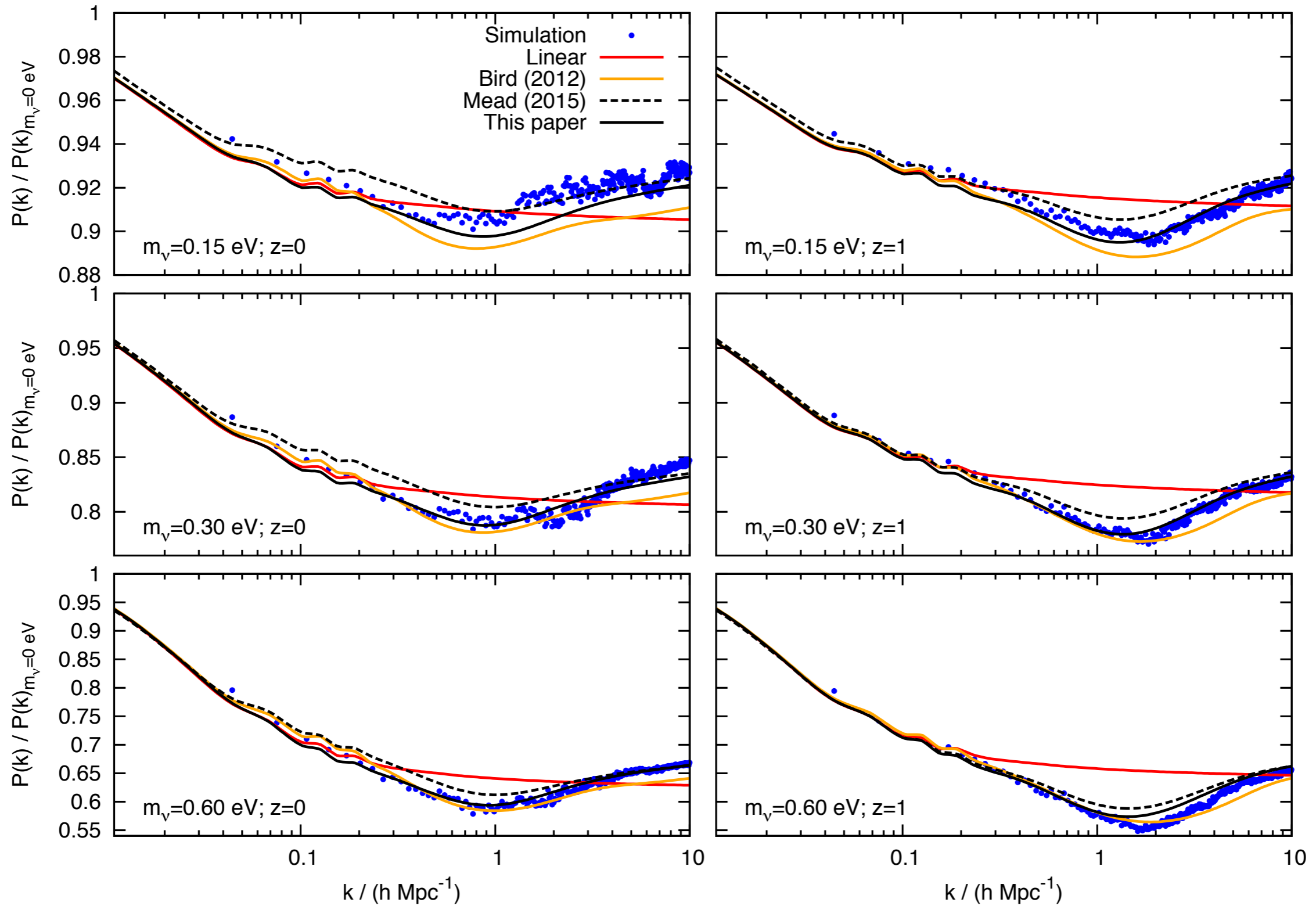


Dark energy

$$w = w_0 + (1 - a)w_a$$



Massive neutrinos



Summary

- Matter power spectra accurate to 5% across a wide range of cosmological models (massive ν , dark energy)
- Publications
 - <http://mnras.oxfordjournals.org/content/454/2/1958.full.pdf>
 - <http://mnras.oxfordjournals.org/content/459/2/1468.full.pdf>
- Code available:
 - <https://github.com/alexander-mead/HMcode>
- Integrated into CAMB (halofit_ppf.f90 module)
- Relatively easy to add standard model extensions

Spherical collapse model

- 5% accuracy can be improved on if we work only with the power spectrum 'response'
- Consider models with fixed linear spectrum shape and amplitude, but that could differ via their growth history (e.g., dark energy)
- Look at differences in spherical-collapse model predictions.

Spherical-collapse model

$$\ddot{\delta} + 2H\dot{\delta} - \frac{4}{3} \frac{\dot{\delta}^2}{1 + \delta} = 4\pi G \bar{\rho} \delta (1 + \delta)$$



Spherical-collapse model

$$\ddot{\delta} + 2H\dot{\delta} - \frac{4}{3} \frac{\dot{\delta}^2}{1 + \delta} = 4\pi G \bar{\rho} \delta (1 + \delta)$$



Spherical-collapse model

$$\ddot{\delta} + 2H\dot{\delta} - \frac{4}{3} \frac{\dot{\delta}^2}{1 + \delta} = 4\pi G \bar{\rho} \delta (1 + \delta)$$

$$\ddot{\delta} + 2H\dot{\delta} = 4\pi G \bar{\rho} \delta$$



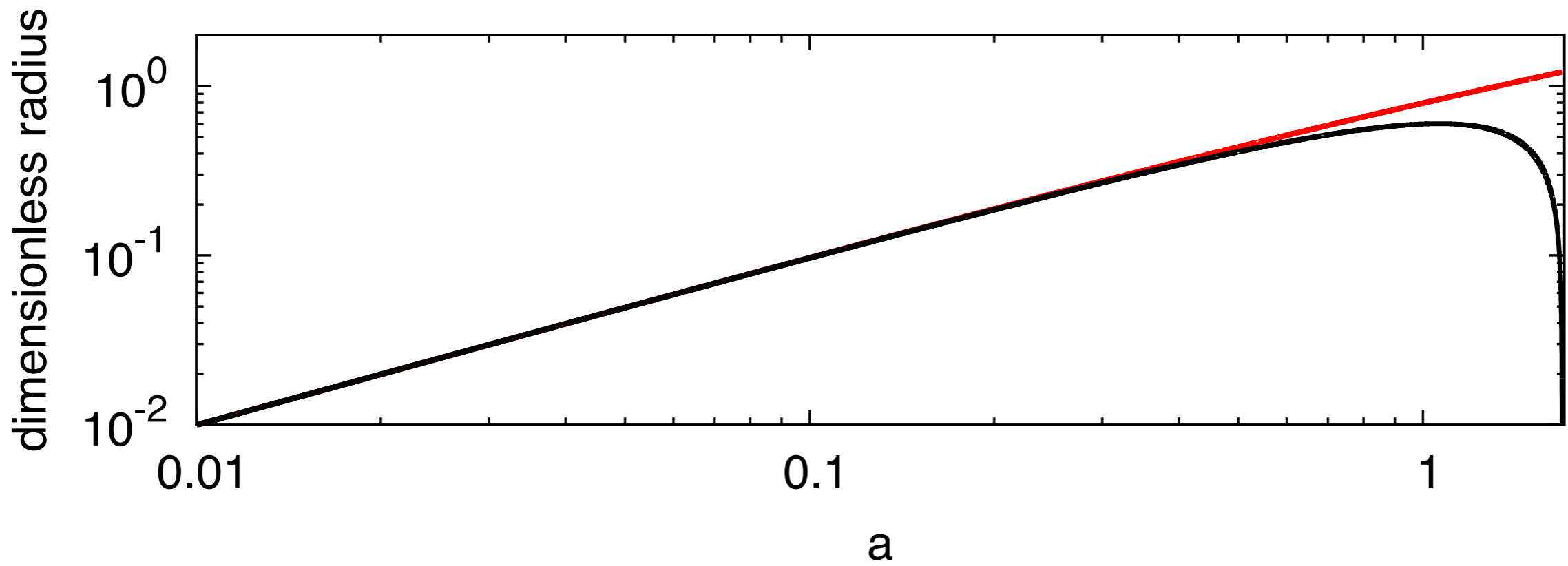
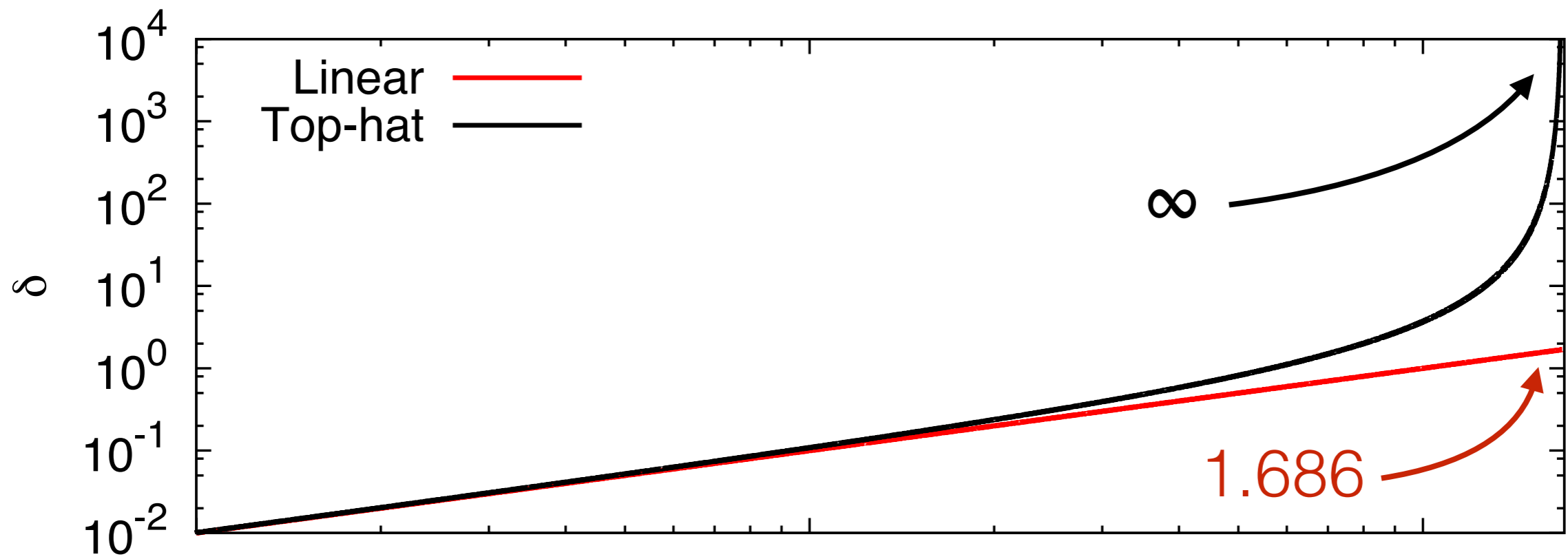
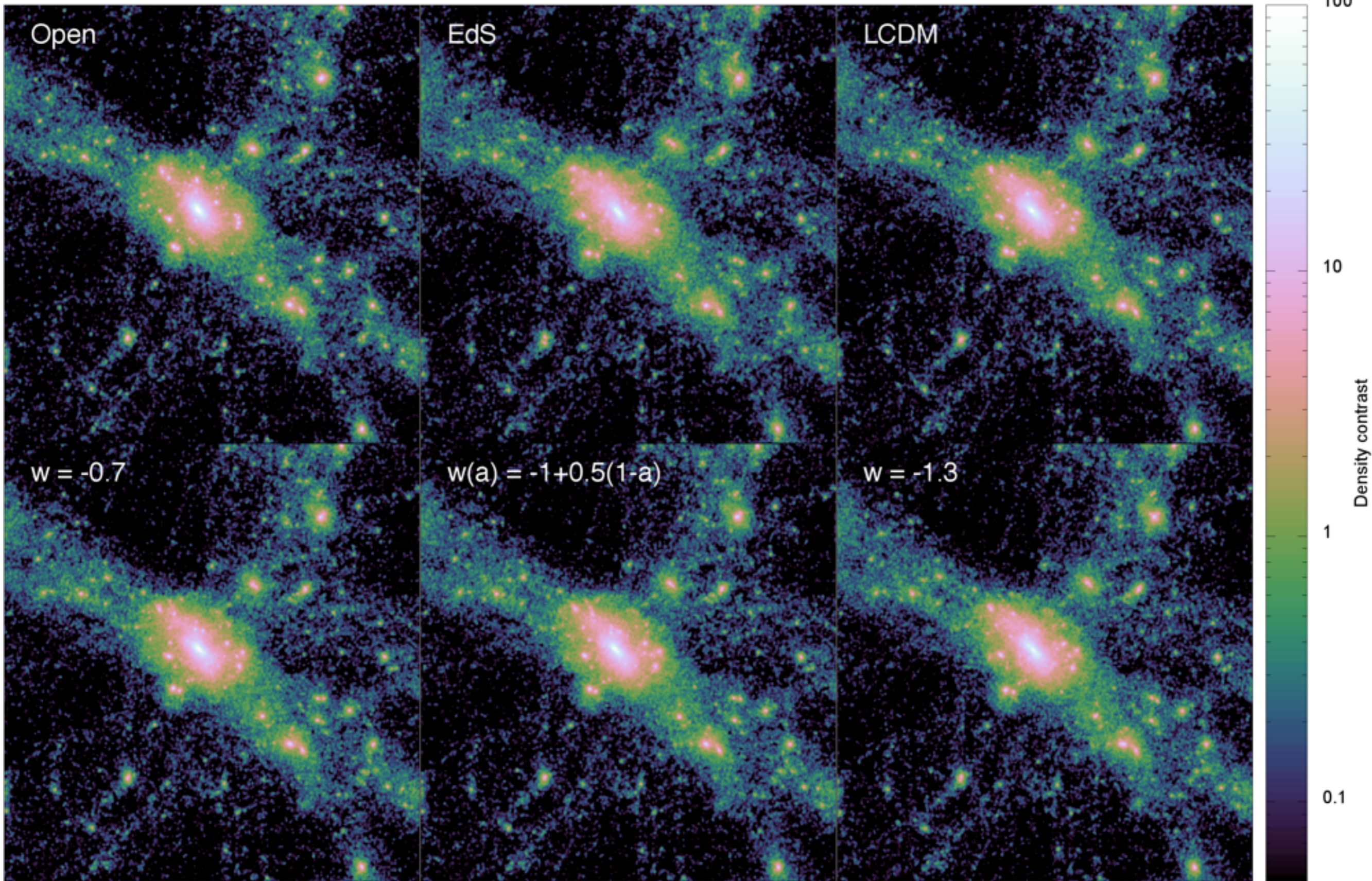
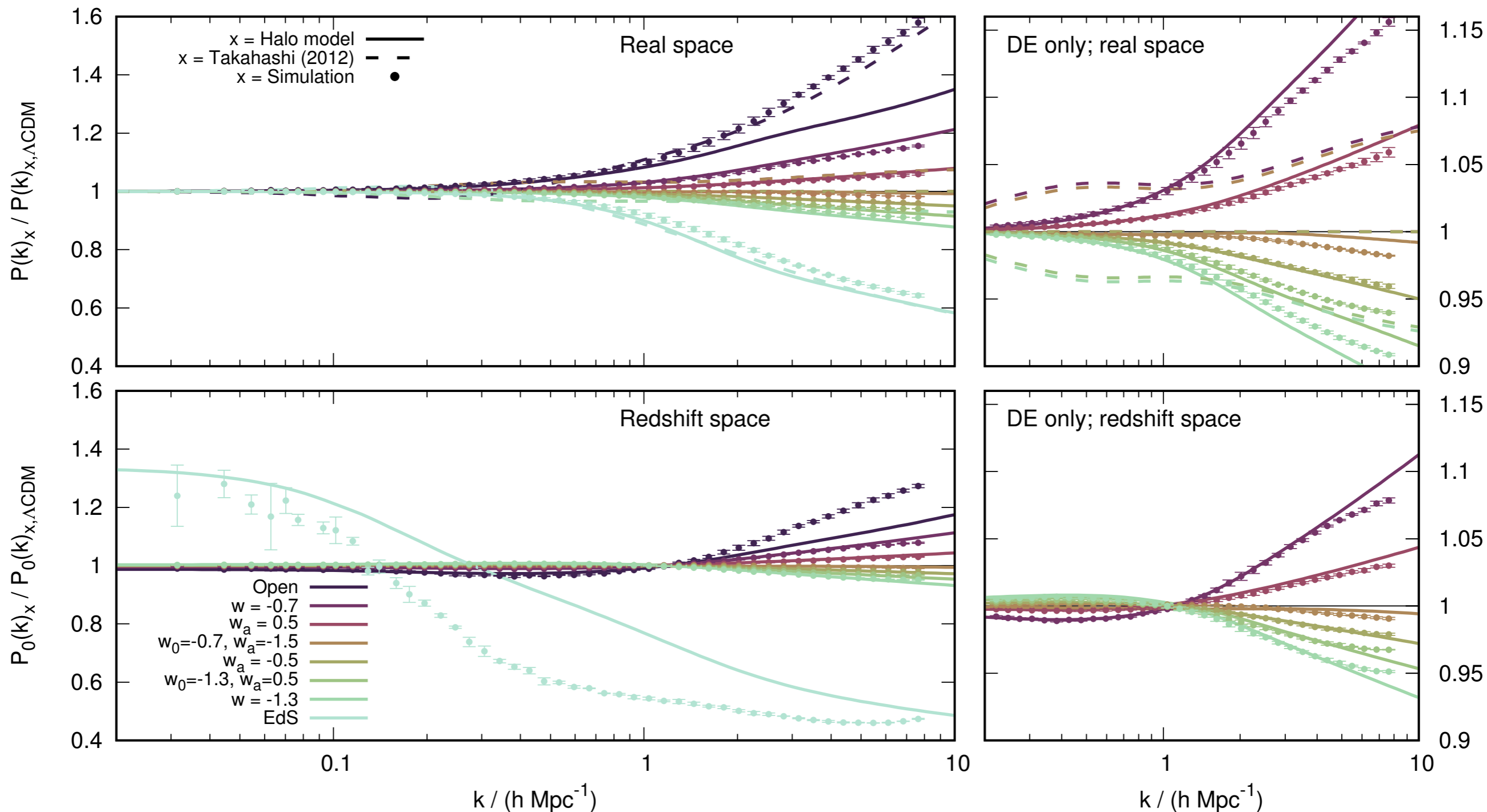


Table 1. Cosmological parameters of the simulations used in this paper. Dynamical dark energy is parameterised via $w(a) = w_0 + (1 - a)w_a$ and is taken to be spatially homogeneous and thus only affects the background expansion. All simulations use 512^3 particles in cubes of size $L = 200 h^{-1} \text{Mpc}$, and start from initial conditions with identical mode phases, but with initial amplitudes adjusted to ensure $\sigma_8 = 0.8$ at $z = 0$. The shape of the linear spectrum used to generate the initial conditions is identical in each case, and was generated using CAMB (Lewis et al. 2000) with cosmological parameters $\Omega_m = 0.3$, $\Omega_w = 1 - \Omega_m$, $\Omega_b = 0.05$, $h = 0.7$, $n_s = 0.96$ and $w = -1$. For each cosmology, I ran 3 different realisations of initial conditions. Also shown are the spherical-model parameters δ_c and Δ_v from a numerical calculation.

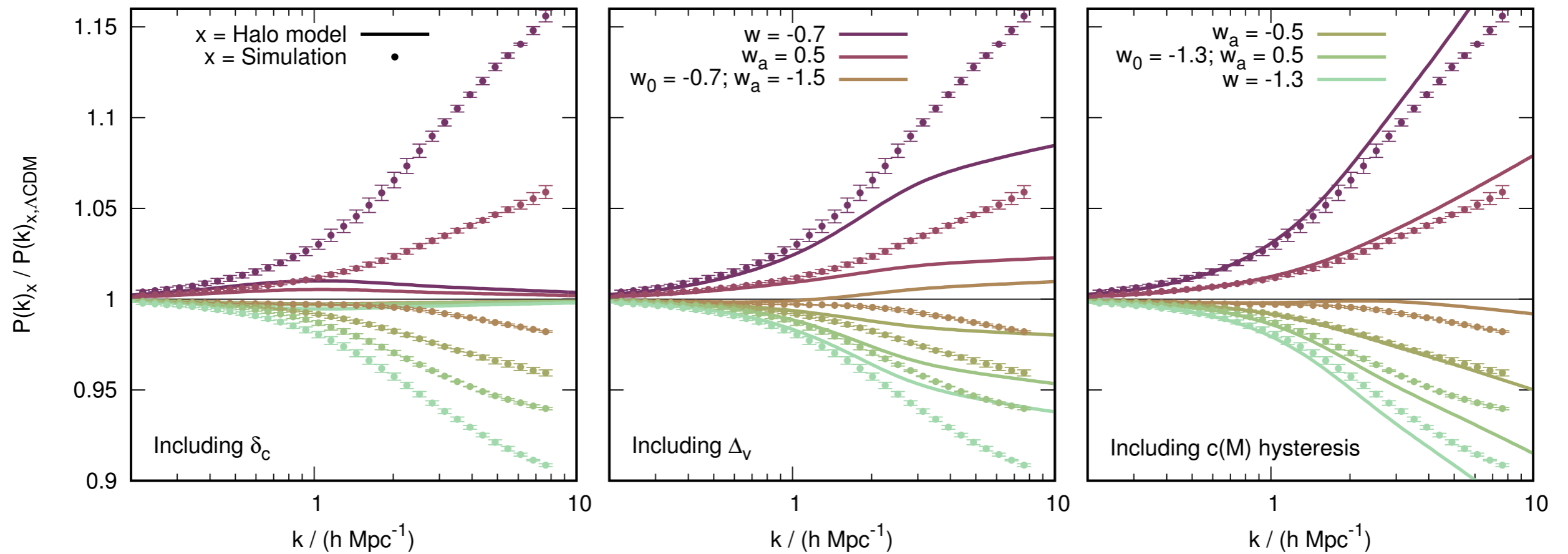
Cosmology	Ω_m	Ω_w	w_0	w_a	δ_c	Δ_v
Λ CDM	0.3	0.7	-1	0	1.6755	310.1
EdS	1.0	0.0	-	-	1.6866	177.7
Open	0.3	0.0	-	-	1.6513	402.0
DE1	0.3	0.7	-0.7	0	1.6695	342.7
DE2	0.3	0.7	-1.3	0	1.6787	282.4
DE3	0.3	0.7	-1	0.5	1.6724	318.5
DE4	0.3	0.7	-1	-0.5	1.6773	301.6
DE5	0.3	0.7	-0.7	-1.5	1.6774	313.3
DE6	0.3	0.7	-1.3	0.5	1.6771	290.1



Spherical collapse model



Spherical collapse model



Summary

- Matter 'power spectrum response' accurate to 1-2% for $k < 5h \text{ Mpc}^{-1}$ (Euclid or LSST lensing accuracy)
- No fitted parameters
- Publication:
 - <http://mnras.oxfordjournals.org/content/early/2016/09/13/mnras.stw2312.full.pdf>
- Code available:
 - <https://github.com/alexander-mead/collapse>
- Papers also contain accurate fitting functions for δ_c and Δ_v for a wide range of dark energy models